# Interacting Particle Systems: Almost sure uniqueness, pathwise duality, and the mean-field limit 

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## Lecture 2: Pathwise duality

## Poisson construction

$S$ a finite set called local state space,
$\Lambda$ a countable set called lattice,
$\mathcal{G}$ a countable set of local maps $m: S^{\wedge} \rightarrow S^{\wedge}$, $\left(r_{m}\right)_{m \in \mathcal{G}}$ nonnegative rates,
$\pi$ Poisson subset of $\mathcal{G} \times \mathbb{R}$ with intensity $r_{m} \mathrm{~d} t$, $\left(\mathbb{X}_{s, u}\right)_{s \leq u}$ stochastic flow constructed from $\pi$.

Theorem Assuming (SUM), the process $\left(X_{t}\right)_{t \geq 0}$ defined as

$$
X_{t}:=\mathbb{X}_{s, s+t}\left(X_{0}\right) \quad(t \geq 0)
$$

is the interacting particle system with generator

$$
\begin{equation*}
G f(x)=\sum_{m \in \mathcal{G}} r_{m}\{f(m(x))-f(x)\} \tag{GEN}
\end{equation*}
$$

## Ergodicity

Theorem Assume that the constant

$$
K:=\sup _{i \in \Lambda} \sum_{\substack{m \in \mathcal{G} \\ \mathcal{D}(m) \ni i}} r_{m}(|\mathcal{R}(m[i])|-1)
$$

satisfies $K<0$. Then the interacting particle system $X$ has a unique invariant law $\nu$, and the process started in any $X_{0}=x$ satisfies

$$
\mathbb{P}^{x}\left[X_{t} \in \cdot\right] \underset{t \rightarrow \infty}{\Longrightarrow} \nu \quad\left(x \in S^{\wedge}\right)
$$

Proof For each $i \in \Lambda, u \in \mathbb{R}$, there exists a random $s>-\infty$ such that $\mathbb{X}_{s, u}[i](x)$ is constant as a function of $x$. In other words, if $u-s$ is large, then $\mathbb{X}_{s, u}[i]$ "forgets" the initial state.

## The backwards in time picture

We may view $\mathbb{X}_{s, u}[i]$ as a map from $S^{\mathcal{R}\left(\mathbb{X}_{s, u}[i]\right)}$ to $S$.
Observation For fixed $u \in \mathbb{R}$ and $i \in \Lambda$, the process

$$
\left(\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right), \mathbb{X}_{u-t, u}[i]\right)_{t \geq 0}
$$

is a Markov process with countable state space consisting of all pairs $(\Delta, \mathbb{X})$ where $\Delta \subset \Lambda$ is finite and $\mathbb{X}: S^{\Delta} \rightarrow S$.

## The dual picture

One has

$$
\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right) \subset \zeta_{u-t}^{u}(\{i\}) \quad \text { with } \quad \mathbb{E}\left[\left|\zeta_{u-t}^{u}(\{i\})\right|\right] \leq e^{K t}
$$

Therefore $K<0$ implies that for $t$ large enough, $\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right)=\emptyset$ and hence $\mathbb{X}_{u-t, u}[i]$ is constant, which means that the interacting particle system "forgets" its initial state and is ergodic.

On the other hand, for $K>0$, it may happen that the system is nonergodic and there are multiple invariant laws.
Much information is contained in the "backward in time" Markov process

$$
\left(\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right), \mathbb{X}_{u-t, u}[i]\right)_{t \geq 0}
$$

It depends on the details of the system under consideration how tractable this backward process is.

## Additive systems

Let $\underline{0}$ denote the configuration that is identically zero.
A map $m:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$ is additive if
(i) $m(\underline{0})=\underline{0}$,
(ii) $m(x \vee y)=m(x) \vee m(y) \quad\left(x, y \in\{0,1\}^{\wedge}\right)$.

Lemma If all maps in $\mathcal{G}$ are additive, then

$$
\mathcal{R}\left(\mathbb{X}_{s, u}[i]\right)=\zeta_{s}^{u}(\{i\}) \quad \text { and } \quad \mathbb{X}_{s, u}[i](x)=\bigvee_{j \in \zeta_{s}^{u}(\{i\})} x(j)
$$

## The contact process


$\mathbb{X}_{s, u}(x)(i)=1$ if and only if $x(j)=1$ for at least one $j \in \zeta_{s}^{u}(\{i\})$.

## Self-duality of the contact process

For the contact process, setting

$$
Y_{t}(i):= \begin{cases}1 & \text { if } i \in \zeta_{u-t}^{u}(A) \\ 0 & \text { otherwise }\end{cases}
$$

defines a "dual" contact process, albeit with left-continuous paths.
Generalising our construction, starting from a graphical representation $\pi$, we can define two stochastic flows $\left(\mathbb{X}_{s, u}^{ \pm}\right)_{s \leq u}$, where:

$$
\begin{aligned}
& \mathbb{X}_{s, u}^{-}:=\text {the concatenation of all } m \text { with }(m, t) \in \pi, t \in[s, u), \\
& \mathbb{X}_{s, u}^{+}:=\text {the concatenation of all } m \text { with }(m, t) \in \pi, t \in(s, u] .
\end{aligned}
$$

In particular, $\mathbb{X}_{s, u}^{+}=\mathbb{X}_{s, u}$ as before.

## Additive systems duality

Let $\psi:\{0,1\}^{\wedge} \times\{0,1\}^{\wedge} \rightarrow\{0,1\}$ be the duality function

$$
\Psi(x, y):=\bigvee_{i \in \Lambda} x(i) y(i) \quad\left(x, y \in\{0,1\}^{\wedge}\right)
$$

Lemma For each additive local map $m:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$, there exists a unique "dual" map $\hat{m}:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$ such that

$$
\Psi(m(x), y)=\Psi(x, \hat{m}(y)) \quad\left(x, y \in\{0,1\}^{\wedge}\right),
$$

and $\hat{m}$ is also an additive local map.
Pathwise duality Let $\hat{\pi}:=\{(\hat{m},-t):(m, t) \in \pi\}$.
Construct $\left(\mathbb{X}_{s, u}^{ \pm}\right)_{s \leq u}$ from $\pi$ and $\left(\mathbb{Y}_{s, u}^{ \pm}\right)_{s \leq u}$ from $\hat{\pi}$. Then:

$$
\Psi\left(\mathbb{X}_{s, u}^{ \pm}(x), y\right)=\Psi\left(x, \mathbb{Y}_{-u,-s}^{\mp}(y)\right) \quad\left(x, y \in\{0,1\}^{\wedge}\right) .
$$

## Dual local maps

Dual maps can be found using the recipe: reverse the arrows, keep the blocking symbols.


## Cancellative systems

Let $(x \oplus y)(i):=x(i)+y(i) \bmod (2)$.
A map $m:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$ is cancellative if
(i) $m(\underline{0})=\underline{0}$,
(ii) $m(x \oplus y)=m(x) \oplus m(y) \quad\left(x, y \in\{0,1\}^{\wedge}\right)$.

Let $|y|:=\sum_{i \in \Lambda} y(i)$ and let $\Psi$ be the duality function

$$
\Psi(x, y):=\bigoplus_{i<\Lambda} x(i) y(i) \quad\left(x, y \in\{0,1\}^{\wedge},|y|<\infty\right)
$$

Lemma For each cancellative local map $m:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$, there exists a unique "dual" map $\hat{m}:\{0,1\}^{\wedge} \rightarrow\{0,1\}^{\wedge}$ such that

$$
\Psi(m(x), y)=\Psi(x, \hat{m}(y)) \quad\left(x, y \in\{0,1\}^{\wedge},|y|<\infty\right)
$$

and $\hat{m}$ is also a cancellative local map.

## Cancellative duality

Pathwise duality Let $\hat{\pi}:=\{(\hat{m},-t):(m, t) \in \pi\}$.
Construct $\left(\mathbb{X}_{s, u}^{ \pm}\right)_{s \leq u}$ from $\pi$ and $\left(\mathbb{Y}_{s, u}^{ \pm}\right)_{s \leq u}$ from $\hat{\pi}$. Then:

$$
\Psi\left(\mathbb{X}_{s, u}^{ \pm}(x), y\right)=\Psi\left(x, \mathbb{Y}_{-u,-s}^{\mp}(y)\right) \quad\left(x, y \in\{0,1\}^{\wedge},|y|<\infty\right)
$$

Now the "backward in time" Markov process

$$
\left(\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right), \mathbb{X}_{u-t, u}[i]\right)_{t \geq 0}
$$

has the simple form

$$
1_{\mathcal{R}\left(\mathbb{X}_{s, u}[i]\right)}=\mathbb{Y}_{-u,-s}^{-}\left(1_{\{i\}}\right)
$$

and

$$
\mathbb{X}_{s, u}[i](x)=\bigoplus_{i \in \mathcal{R}\left(\mathbb{X}_{s, u}[i]\right)} x(j)
$$

## The annihilating branching process


$\mathbb{X}_{s, u}(x)(i)=1$ if and only if $\left|\left\{j \in \mathcal{R}\left(\mathbb{X}_{s, u}[i]\right): x(j)=1\right\}\right|$ is odd.

## Threshold voter model

For any graph $(\Lambda, \sim), i \in \Lambda$, and $x \in\{0,1\}^{\wedge}$, let

$$
\begin{aligned}
& \operatorname{thresh}_{i}^{0}(x)(j):= \begin{cases}\bigwedge_{k: k \sim i} x(k) & \text { if } j=i, \\
x(j) & \text { otherwise. }\end{cases} \\
& \operatorname{thresh}_{i}^{1}(x)(j):= \begin{cases}\bigvee_{k: k \sim i} x(k) & \text { if } j=i, \\
x(j) & \text { otherwise. }\end{cases}
\end{aligned}
$$

The threshold voter model is defined by the generator

$$
\begin{aligned}
G_{\text {thresh }} f(x):= & \sum_{i \in \Lambda}\left\{f\left(\operatorname{thresh}_{i}^{0}(x)\right)-f(x)\right\} \\
& +\sum_{i \in \Lambda}\left\{f\left(\operatorname{thresh}_{i}^{1}(x)\right)-f(x)\right\} .
\end{aligned}
$$

Note thresh ${ }_{i}^{1}$ is additive but thresh ${ }_{i}^{0}$ is only monotone.

## Threshold voter model

Let $\mathcal{N}_{i}:=\{i\} \cup\{k: k \sim i\}, N:=\left|\mathcal{N}_{i}\right|$, and $\mathcal{D}_{i}:=\left\{\Delta \subset \mathcal{N}_{i}:|\Delta|\right.$ is even $\}$. We define

$$
\mathrm{flip}_{i}^{\Delta}(x)(j):= \begin{cases}\bigoplus_{k \in \Delta} x(k) & \text { if } j=i \\ x(j) & \text { otherwise }\end{cases}
$$

Then

$$
G_{\text {thresh }} f(x)=2^{-N+2} \sum_{i \in \Lambda} \sum_{\Delta \subset \mathcal{D}_{i}}\left\{f\left(\text { flip }_{i}^{\Delta}(x)\right)-f(x)\right\} .
$$

Note $\mathrm{flip}_{i}^{\Delta}$ is cancellative but not monotone.

## Threshold voter model

Trivially, $\delta_{\underline{0}}$ and $\delta_{\underline{1}}$ are invariant laws.
Def a coexisting probability law gives zero probability to $\{\underline{0}, \underline{1}\}$.
Complete convergence [Handjani 1999] The threshold voter model on $\mathbb{Z}^{d}$ with $d \geq 2$ has a unique coexisting invariant law $\nu$. Setting

$$
\begin{aligned}
& p_{0}:=\mathbb{P}\left[X_{t}=\underline{0} \text { for some } t<\infty\right], \\
& p_{1}:=\mathbb{P}\left[X_{t}=\underline{1} \text { for some } t<\infty\right]
\end{aligned}
$$

the process starting in a general initial law satisfies

$$
\mathbb{P}\left[X_{t} \in \cdot\right] \underset{t \rightarrow \infty}{\Longrightarrow} p_{0} \delta_{\underline{0}}+p_{1} \delta_{\underline{1}}+\left(1-p_{0}-p_{1}\right) \nu
$$

## Threshold voter model

Cancellative duality gives

$$
\Psi\left(X_{t}, y\right)=\Psi\left(\mathbb{X}_{0, t}^{+}\left(X_{0}\right), y\right)=\Psi\left(\mathbb{X}_{0, s}^{+}\left(X_{0}\right), \mathbb{Y}_{-t,-s}^{-}(y)\right)
$$

Denote the survival probability of the dual by:

$$
\rho(y):=\lim _{t \rightarrow \infty} \mathbb{P}\left[\mathbb{Y}_{-t, 0}^{-}(y) \neq \underline{0}\right]
$$

The proof consists of showing that

$$
\lim _{t \rightarrow \infty} \mathbb{E}\left[\Psi\left(X_{t}, y\right)\right]=p_{0} \Psi(\underline{0}, y)+p_{1} \Psi(\underline{1}, y)+\frac{1}{2}\left(1-p_{0}-p_{1}\right) \rho(y)
$$

Conditional on non-absorption of $\mathbb{Y}_{-t,-s}^{-}(y)$ in $\underline{0}$, the 1's spread over all of space,
Conditional on non-absorption of $\mathbb{X}_{0, s}^{+}\left(X_{0}\right)$ in $\underline{0}$ or $\underline{1}$, there are many 0 's close to 1 's.

## Generalisations of additive and cancellative duality

Remark More generally, one can look at duality functions of the form

$$
\Psi(x, y):=\sum_{i \in \Lambda} \psi(x(i), y(i))
$$

where $\psi: S \times S \rightarrow T$ is a "local" duality function taking values in a commutative monoid $(T,+)$, and $\sum_{i \in \Lambda}$ is the sum in $(T,+)$.
$S=T=\{0,1\}$ and $x+y:=x \vee y$ gives additive duality.
$S=T=\{0,1\}$ and $x+y:=x \oplus y$ gives cancellative duality. In both cases $\psi(x, y)=x \cdot y$.

For local state spaces with two elements this seem to be the only two "reasonable" choices but already for $|S|=3$ there are many more possibilities. See [Latz \& S '21].

## The cooperative contact process



A particle system using the maps cob and dth.

## The cooperative contact process

The cooperative branching map is monotone but not additive.

$$
\operatorname{cob}_{i_{1} i_{2} i_{3}}(x)(j):= \begin{cases}x\left(i_{1}\right) \vee\left(x\left(i_{2}\right) \wedge x\left(i_{3}\right)\right) & \text { if } j=i_{1} \\ x(j) & \text { otherwise }\end{cases}
$$

We can write

$$
\mathbb{X}_{s, u}[i](x)=\bigvee_{\Delta \in \mathcal{Z}_{s, u}(i)} \bigwedge_{j \in \Delta} x(j)
$$

where $\mathcal{Z}_{s, u}(i)$ is the set of "minimal configurations" $\Delta$ which need to be 1 in order for $\mathbb{X}_{s, u}[i](x)$ to be 1 .

## Minimal configurations



## Minimal configurations



## Minimal configurations



## Minimal configurations



## Monotone systems duality

Recall the "backward in time" Markov process

$$
\left(\mathcal{R}\left(\mathbb{X}_{u-t, u}[i]\right), \mathbb{X}_{u-t, u}[i]\right)_{t \geq 0}
$$

For monotone systems,

$$
\begin{aligned}
\mathcal{R}\left(\mathbb{X}_{s, u}[i]\right) & =\bigcup_{\Delta \in \mathcal{Z}_{s, u}(i)} \Delta, \\
\mathbb{X}_{s, u}[i](x) & =\bigvee_{\Delta \in \mathcal{Z}_{s, u}(i)} \bigwedge_{j \in \Delta} x(j) .
\end{aligned}
$$

These sort of dual processes have first been used by Gray (1986).
In general, these sort of dual processes are hard to control.
In the next lecture, we will study "backward in time" processes in the mean-field limit.

