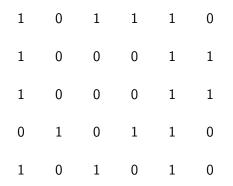
### On rebellious voter models

Jan M. Swart

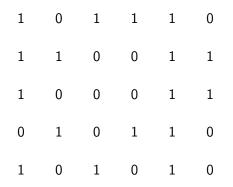
#### Vienna, June 18, 2013 joint with Anja Sturm and Karel Vrbenský

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#### 0's and 1's represent two closely related species.

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#### The 0's and 1's evolve in a Markovian way.

1	0	1	1	1	0
1	1	0	0	1	1
1	0	0	0	1	1
0	0	0	1	1	0
1	0	1	0	1	0

#### The 0's and 1's evolve in a Markovian way.

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1	0	1	1	0	0
1	1	0	0	1	1
1	0	0	0	1	1
0	0	0	1	1	0
1	0	1	0	1	0

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1	0	1	1	0	0
1	1	0	0	1	1
1	0	0	0	1	1
0	0	1	1	1	0
1	0	1	0	1	0

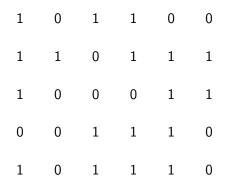
#### The 0's and 1's evolve in a Markovian way.

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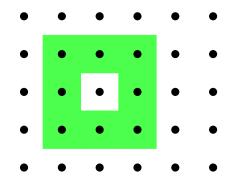
1	0	1	1	0	0
1	1	0	0	1	1
1	0	0	0	1	1
0	0	1	1	1	0
1	0	1	1	1	0

#### The 0's and 1's evolve in a Markovian way.

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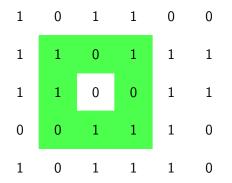


#### The 0's and 1's evolve in a Markovian way.



$$\mathcal{N}_i := \{ j \in \mathbb{Z}^d : 0 < \|i - j\|_{\infty} \le R \} \text{ neighborhood of a site.}$$
  
(Here  $R = 1, d = 2$ ).

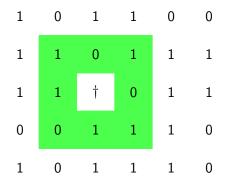
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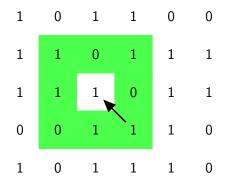
 $f_0 = 3/8$ ,  $f_1 = 5/8$  local frequencies of types 0, 1.

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With rate with rate  $f_0 + \alpha_{01}f_1$  an organism of type 0 dies...



... and is replaced by a random type from the neighborhood.

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**Neuhauser & Pacala (1999):** Markov process in the space  $\{0,1\}^{\mathbb{Z}^d}$  of spin configurations  $x = (x(i))_{i \in \mathbb{Z}^d}$ , where spin x(i) flips:

$$0 \mapsto 1 \text{ with rate } f_1(f_0 + \alpha_{01}f_1),$$
  
 
$$1 \mapsto 0 \text{ with rate } f_0(f_1 + \alpha_{10}f_0),$$

#### with

$$f_{ au}(i):=rac{\#\{j\in\mathcal{N}_i:x(j)= au\}}{\#\mathcal{N}_i}\quad\mathcal{N}_i:=\{j:0<\|i-j\|_\infty\leq R\}.$$

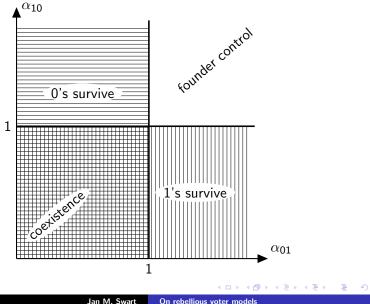
the local frequency of type  $\tau = 0, 1$ .

**Interpretation:** Interspecific competition rates  $\alpha_{01}, \alpha_{10}$ . Organism of type 0 dies with rate  $f_0 + \alpha_{01}f_1$  and is replaced by type sampled at random from distance  $\leq R$ .

By definition, type 0 *survives* if starting from a single organism of type 0 and all other organisms of type 1, there is a positive probability that the organisms of type 0 never die out.

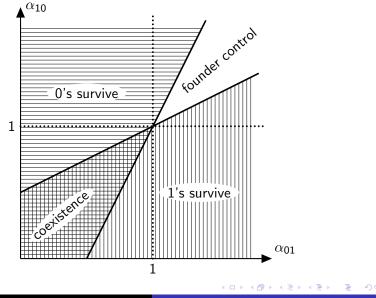
By definition, one has *coexistence* if there exists an invariant law concentrated on states where both types are present.

# Mean field model



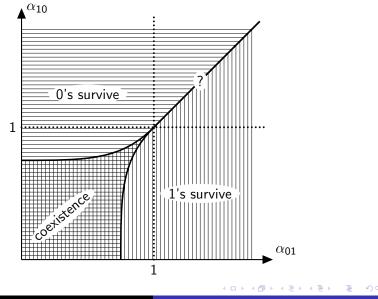
Jan M. Swart

# Dimension $d \ge 3$



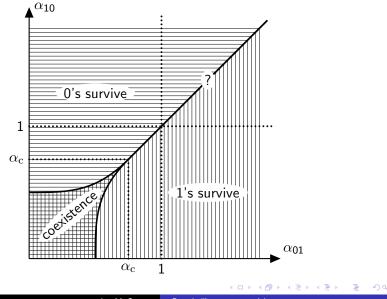
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# Dimension d = 2



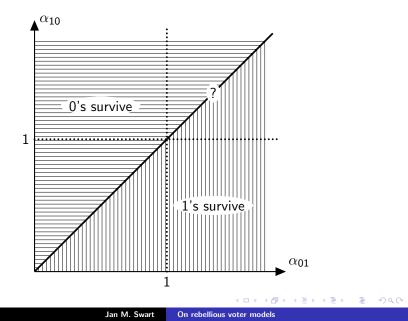
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## Dimension d = 1, range $R \ge 2$



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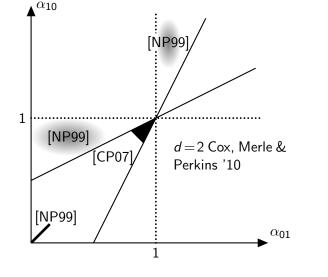
# Dimension $d = \overline{1}$ , range R = 1



**Conjecture** There exists a critical dimension  $d_c \cong 4/3$  such that in dimensions  $d < d_c$ , two species must be sufficiently different to be able to coexist, but in dimensions  $d > d_c$ , any difference, no matter how small, suffices.

(3)

### Rigorous results for $d \ge 3$

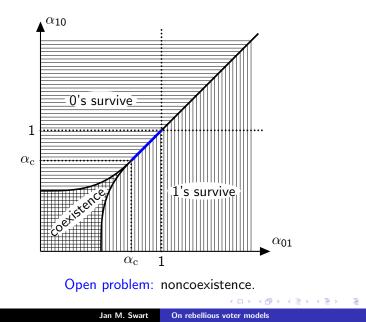


[NP99]=Neuhauser & Pacala '99, [CP07]=Cox & Perkins, '07.

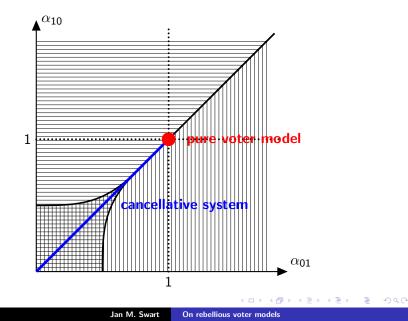
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On rebellious voter models

# Open problem for d = 1



# Special models



#### Pure voter model

If started with finitely many organisms of type 1, then number of 1's is a martingale. Consequence: 1's (and likewise 0's) die out.

Dual to coalescing random walks. Consequence: coexistence in transient dimensions  $d \ge 3$ , clustering in d = 1, 2.

The symmetric case  $\alpha_{01} = \alpha_{10} = \alpha \le 1$  is a *cancellative system*. There is a *dual process* Y such that

$$\mathbb{P}\big[|X_t Y_0| \text{ is odd}\big] = \mathbb{P}\big[|X_0 Y_t| \text{ is odd}\big] \qquad (t \ge 0)$$

whenever X and Y are independent. Here

$$|x| := \sum_i x(i)$$
 and  $xy(i) := x(i)y(i)$ .

Equip  $\{0,1\}$  with the usual product and with addition modulo 2, denoted as  $\oplus$ . Then  $\{0,1\}$  is a *finite field*. We may view  $\{0,1\}^{\mathbb{Z}^d}$  (equipped with  $\oplus$ ) as a *linear space* over  $\{0,1\}$ .

A cancellative system  $X = (X_t)_{t \ge 0}$  is a linear system w.r.t. to the finite field  $\{0, 1\}$ , that evolves as

$$x \mapsto x \oplus Ax$$
 with rate  $r(A) \ge 0$ ,

where

$$Ax(i) := \bigoplus_{j \in \mathbb{Z}^d} A(i,j)x(j)$$

with A(i,j) = 1 for finitely many i, j and A(i, j) = 0 otherwise.

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**Example** For  $k \in \mathbb{Z}$ , define:

$$egin{aligned} & A_k(k-1,k) := 1, \quad A_k(k,k) := 1, \quad A_k(i,j) := 0 \ ext{otherwise}, \ & A_k'(k-2,k) := 1, \quad A_k'(k-1,k) := 1, \quad A_k'(i,j) := 0 \ ext{otherwise}. \end{aligned}$$

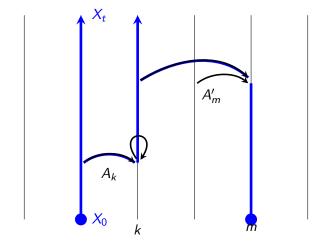
Set  $r(A_k) := \alpha$ ,  $r(A'_k) := 1 = \alpha$ , and r(A) := 0 for all other A. This yields *one-sided rebellious voter model* where x(k) flips

 $0 \leftrightarrow 1$  with rate  $\alpha \mathbb{1}_{\{x(k-1) \neq x(k)\}} + (1-\alpha) \mathbb{1}_{\{x(k-2) \neq x(k-1)\}}$ .

Draw space horizontally, time vertically.

If the local map A applies at time t, draw an arrow from (i, t) to (j, t)whenever A(i, j) = 1.

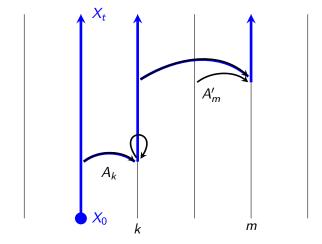
## Graphical representation



 $X_t(i) = 1$  iff there is a odd number of paths from  $X_0$  to (i, t).

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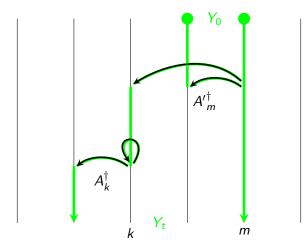
## Graphical representation



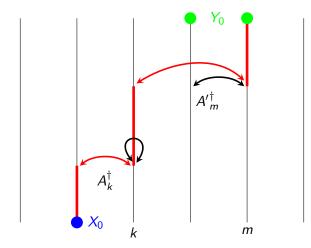
 $X_t(i) = 1$  iff there is a odd number of paths from  $X_0$  to (i, t).

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## Graphical representation of dual model



Time runs backwards and all arrows are reversed.



 $|X_0Y_t|$  is odd  $\Leftrightarrow \#$  paths from  $X_0$  to  $Y_0$  is odd  $\Leftrightarrow |X_tY_0|$  is odd.

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Rates of the dual model:

$$r_Y(A^{\dagger})=r_X(A),$$

where  $A^{\dagger}(i, j) = A(j, i)$  denotes the *adjoint* of *A*. **Duality:** 

$$\mathbb{P}\big[|X_t Y_0| \text{ is odd}\big] = \mathbb{P}\big[|X_0 Y_t| \text{ is odd}\big] \qquad (t \ge 0)$$

whenever X and Y are independent, where

$$|x| := \sum_i x(i)$$
 and  $xy(i) := x(i)y(i)$ .

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If X has voter model dynamics:

$$ig(x(k-1),x(k)ig)\mapstoig(x(k-1),x(k-1)ig) \quad ext{with rate } lpha,$$

then Y has annihilating random walk dynamics:

$$(y(k-1), y(k)) \mapsto (y(k-1) \oplus y(k), 0)$$
 with rate  $\alpha 1_{\{y(k)=1\}}$ ,

i.e., a particle at k jumps to k - 1; if there is already a particle at k - 1, the two particles *annihilate*. If X has rebellious dynamics:

$$x(k)\mapsto 1-x(k)$$
 with rate  $(1-lpha)1_{\{x(k-2)
eq x(k-1)\}}$ ,

then Y has annihilating branching dynamics:

$$y \mapsto y \oplus \delta_{k-2} \oplus \delta_{k-1}$$
, with rate  $(1 - \alpha) \mathbb{1}_{\{y(k)=1\}}$ ,

i.e., a particle at k produces two new particles at positions k - 2 and k - 1and these particles *annihilate* with any particles that may already be present. Similarly, the Neuhauser-Pacala model is dual to a system of *branching-annihilating* particles, where: particles jump with rate  $\alpha$  to a new place at distance  $\leq R$ , and each particle produces two new particles at distances  $\leq R$  with rates proportional to  $1 - \alpha$ .

Since the number of particles of Y is always increased or decreased by 2, the process is *parity preserving*, i.e.,  $|Y_t|$  is odd  $\Leftrightarrow |Y_0|$  is odd.

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By definition, Y survives if starting from an *even* number of particles, there is a positive probability that the particles never die out.

By definition, Y is *stable* if starting from *one* particle, the process 'modulo translations' is positively recurrent, i.e., the system returns to a state with only one particle infinitely often and spends a positive fraction of its time in such states.

By definition, Y is *persistent* if there exists an invariant law concentrated on nonempty configurations.

**Lemma** X has coexistence  $\Leftrightarrow$  Y survives.

**Proof** Start X in product measure with intensity 1/2 and let  $Y_0 := \delta_i + \delta_j$ . Then

$$\begin{split} \mathbb{P}\big[X_t(i) \neq X_t(j)\big] &= \mathbb{P}\big[|X_t Y_0| \text{ is odd}\big] = \mathbb{P}\big[|X_0 Y_t| \text{ is odd}\big] \\ &= \frac{1}{2}\mathbb{P}\big[Y_t \neq 0\big] \xrightarrow[t \to \infty]{} \frac{1}{2}\mathbb{P}\big[Y_s \neq 0 \ \forall s \ge 0\big]. \end{split}$$

**Lemma** In X both types survive  $\Leftrightarrow$  Y is persistent.

For small  $\alpha$ , coexistence and survival of both types for the Neuhauser-Pacala model proved by showing that dual model survives and is persistent.

Main tool: comparison with oriented percolation.

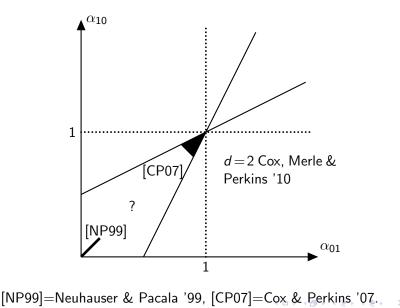
In dimensions  $d \ge 2$ , Cox, Merle and Perkins prove that as  $\alpha_{01}, \alpha_{10} \rightarrow 1$  and  $\alpha_{01} \approx \alpha_{10}$ , rescaled sparse models converge to supercritical super Brownian motion.

Using this, for  $\alpha_{01}, \alpha_{10}$  fixed but very close to *one*, they are able to compare with oriented percolation and prove coexistence.

Intermediate  $\alpha$  still open since monotonicity not proved.

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## Coexistence results for $d \ge 3$



Jan M. Swart

On rebellious voter models

## Interfaces

By definition, a cancellative spin-system X is *type-symmetric* if it treats the types symmetrically, i.e.,  $1 - X_t$  has the same dynamics as  $X_t$ .

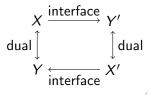
**Lemma** X type-symmetric  $\Leftrightarrow$  dual Y parity preserving.

For one-dimensional, type-symmetric X, setting

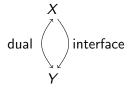
$$Y'_t(i) := 1_{\{X_t(i-rac{1}{2}) \neq X_t(i+rac{1}{2})\}}$$
  $(i \in \mathbb{Z} + rac{1}{2})$ 

defines the *interface model* of X.

**Lemma [Swa13]** Y' is a also a parity preserving cancellative system and



For the rebellious voter model, the dual and interface models coincide:



**Lemma [SS08]** For the rebellious voter model both types survive  $\Leftrightarrow$  the model exhibits coexistence. By definition, X exhibits *interface tightness* if its interface model Y' is stable, i.e., Y' modulo translations is positively recurrent.

Interface tightness means that starting from ... 000000111111..., the system spends a positive fraction of time in such states: the types *cannot invade* each other's territory.

Interface tightness for long-range voter models was proved by Cox and Durrett (1995) under a third moment condition on the infection rates. This was improved to a second moment condition, which is sharp, by Belhaouari, Mountford and Valle (2007). A simpler proof was given by S. & Sturm (2008).

Interface tightness for the Neuhauser-Pacala model with  $R \ge 2$  or for the rebellious voter model is an *open problem*.

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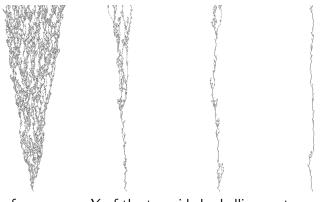
By definition, X exhibits strong interface tightness if for the invariant law of the interface model modulo shifts  $\mathbb{E}[|Y'_{\infty}|] < \infty$ , i.e., the expected number of sites such that  $X_t(i - \frac{1}{2}) \neq X_t(i + \frac{1}{2})$  is finite.

Strong interface tightness is known to hold for voter models and numerically observed for the rebellious voter model with  $\alpha > \alpha_c$ . By contrast, the expected *length* of the interface (distance from left-most one to right-most zero) is known to be infinite for voter models.

**Theorem [Swa13]** Strong interface tightness implies noncoexistence.

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## Numerical simulation

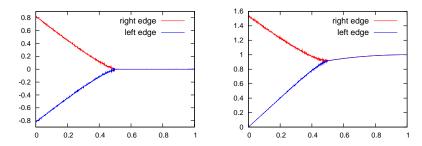


Interface process Y of the two-sided rebellious voter model for  $\alpha = {\rm 0.4, 0.5, 0.51, 0.6.}$ 

## One-sided rebellious interface model



Interface process Y of the one-sided rebellious voter model for  $\alpha = 0.3, 0.5, 0.6.$ 



Edge speeds for the rebellious voter model (left) and its one-sided counterpart (right) [S. & Vrbenský '10].

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Define the survival probability

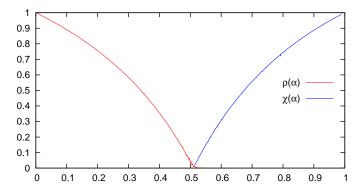
$$\rho(\alpha) := \mathbb{P}^{\delta_0}[X_t \neq 0 \ \forall t \ge 0].$$

• coexistence  $\Leftrightarrow \rho(\alpha) > 0$ .

Define the fraction of time spent with a single interface

$$\chi(\alpha) := \mathbb{P}\big[|Y_{\infty}| = 1\big].$$

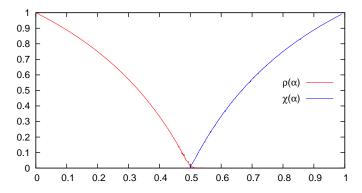
• interface tightness  $\Leftrightarrow \chi(\alpha) > 0$ .



The functions  $\rho$  and  $\chi$  for the two-sided rebelious voter model.

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The functions  $\rho$  and  $\chi$  for the one-sided rebelious voter model.

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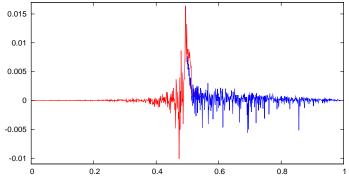
It seems that for the one-sided model, the functions  $\rho$  and  $\chi$  are described by the explicit formulas:

$$ho(lpha) = \mathsf{0} \lor rac{1-2lpha}{1-lpha} \quad ext{and} \quad \chi(lpha) = \mathsf{0} \lor ig(2-rac{1}{lpha}ig).$$

In particular, one has the symmetry  $\rho(1-\alpha) = \chi(\alpha)$  and the critical parameter seems to be given by  $\alpha_c = 1/2$ .

Explanation?

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Differences of  $\rho$  and  $\chi$  with presumed explicit formulas.

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Theoretical physicists believe that

$$ho(lpha)\sim (lpha_{
m c}-lpha)^{eta}$$
 as  $lpha\uparrow lpha_{
m c},$ 

where  $\beta$  is a *critical exponent*.

It has been conjectured by I. Jensen (1994) that  $\beta = 13/14$  and that  $\beta = 1$  [Inui & Tretyakov '98]. More recent estimates are  $\beta \approx 0.92$ ,  $\beta \approx 0.95$  [Hinrichsen '00] [Ódor & Szolnoki '05]. Our formula would imply  $\beta = 1$ .

Let  $\xi^1, \xi^2, \xi^3$  be independent random walks started from  $(\xi_0^1, \xi_0^2, \xi_0^3) = (-1, 0, 1).$ 

Set  $\tau_{ij} := \inf\{t \ge 0 : \xi_t^i = \xi_t^j\}$  and

 $\tau := \tau_{12} \wedge \tau_{23} \wedge \tau_{31}.$ 

Then

$$\mathbb{P}[ au > t] \sim t^{-3/2} \quad ext{as} t o \infty,$$
  
 $\mathbb{E}[ au] = 1 < \infty.$ 

Observed: for small branching rate  $1 - \alpha$ , system Y spends fraction of time of order  $(1 - \alpha)^m$  with 1 + 2m particles.

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