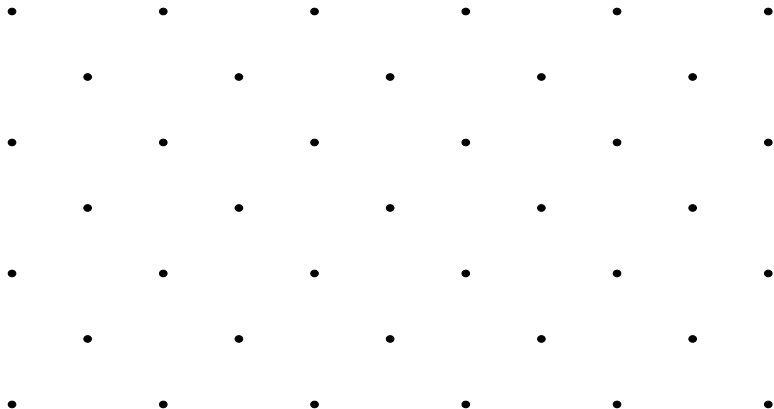


# The Brownian web and net

Jan M. Swart

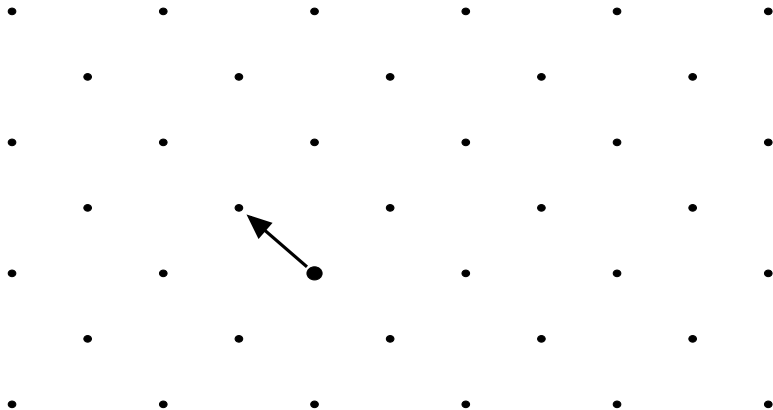
Bedlewo, April 22, 2024

# The even sublattice



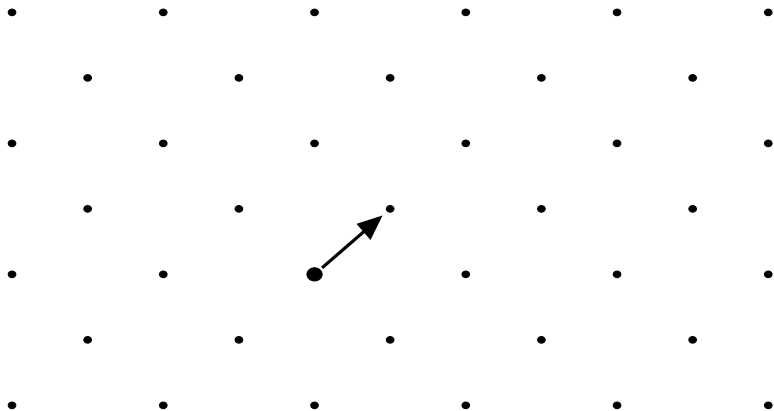
$$\mathbb{Z}_{\text{even}}^2 := \{(x, t) \in \mathbb{Z}^2 : x + t \text{ is even}\}.$$

# Arrow configurations



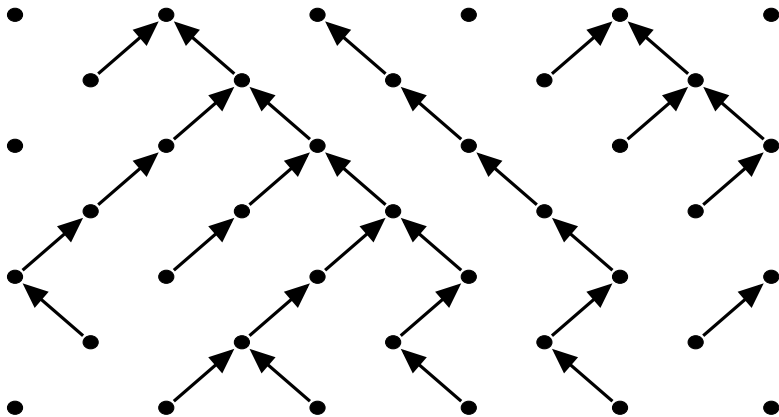
With probability  $\frac{1}{2}$  we draw an arrow to the left...

# Arrow configurations



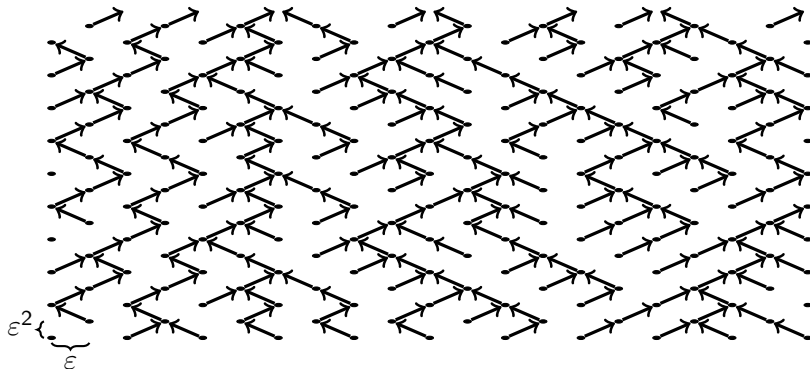
and with probability  $\frac{1}{2}$  we draw it to the right.

# Arrow configurations



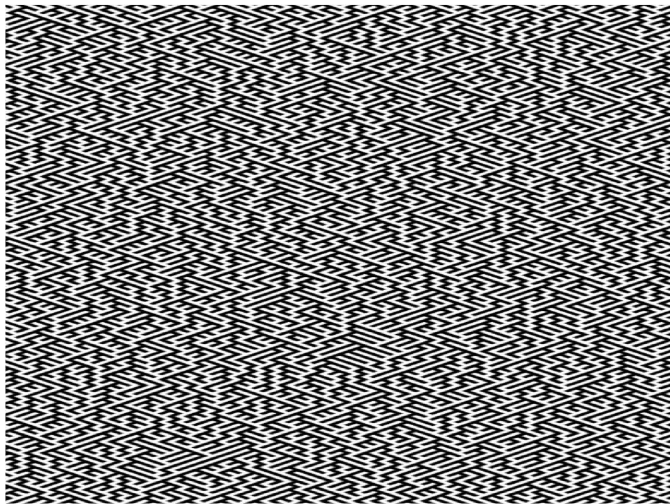
Independently for each space-time point

# Scaling limit



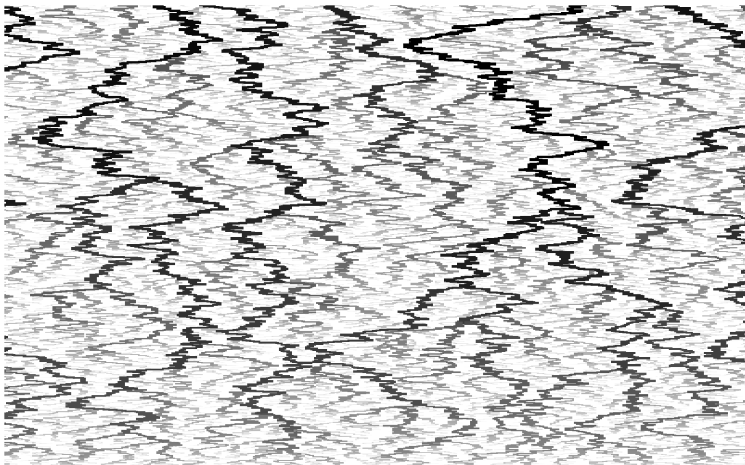
We rescale diffusively, multiplying all spatial distances with  $\varepsilon$  and all temporal distances with  $\varepsilon^2$ .

# Scaling limit



Our aim is to describe the limit as  $\varepsilon \rightarrow 0$ .

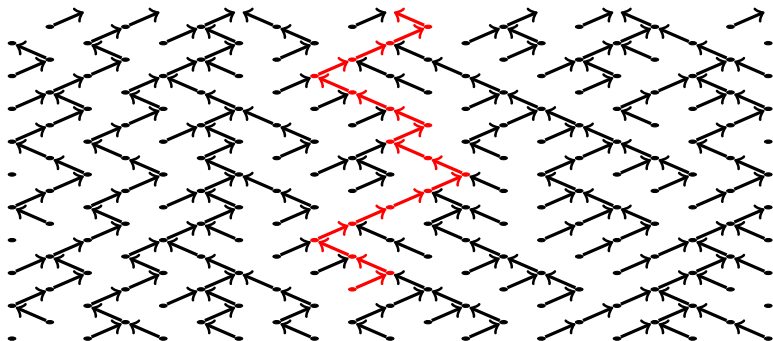
# Scaling limit



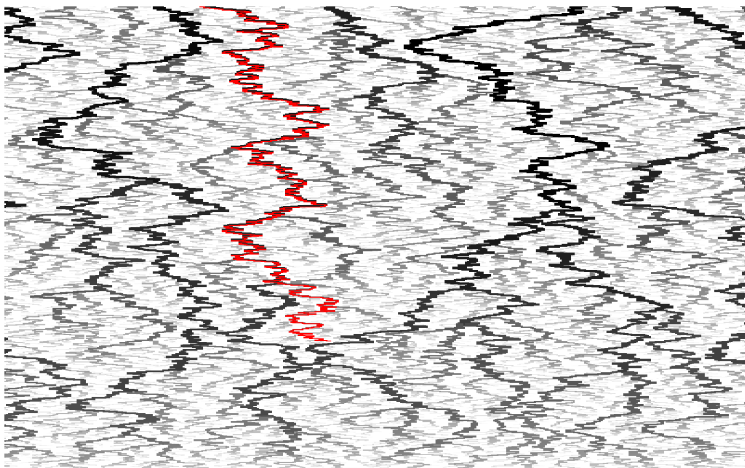
Our aim is to describe the limit as  $\varepsilon \rightarrow 0$ .



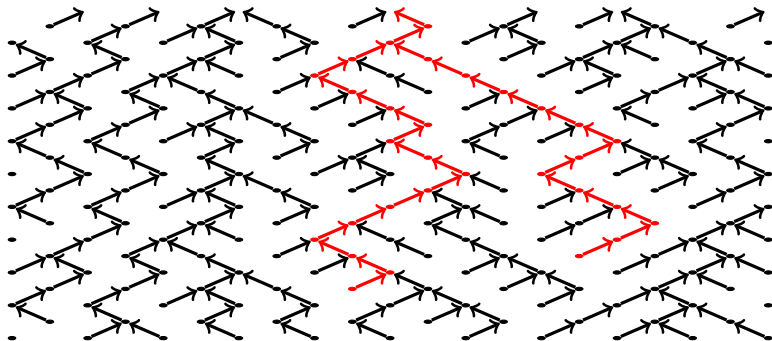
# Paths



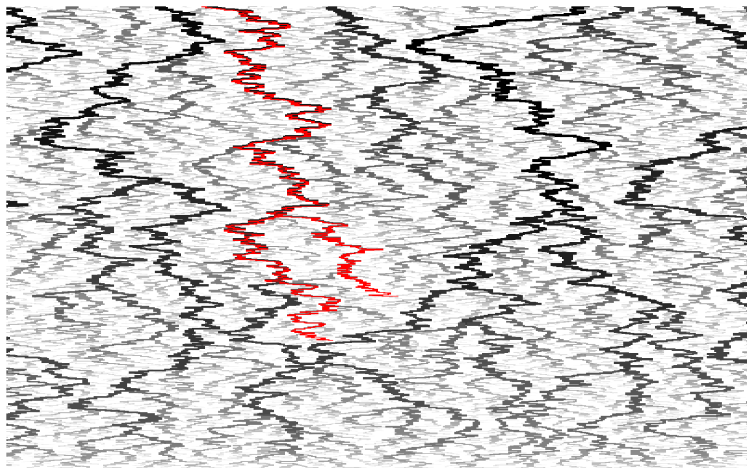
Each space-time point is the starting point of a **random walk path**.



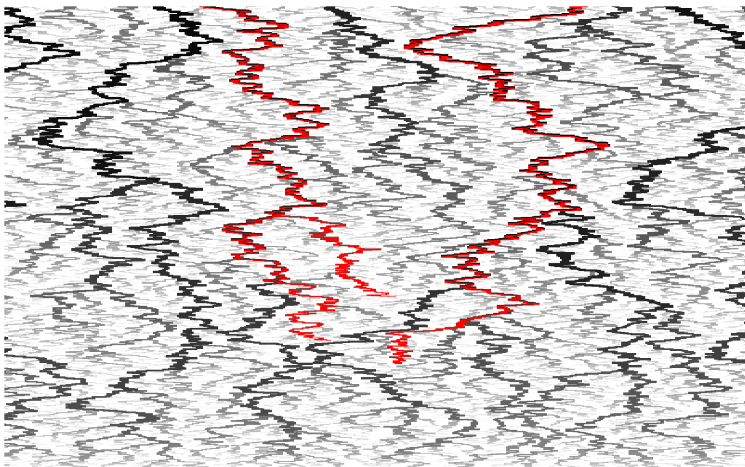
The diffusive scaling limit of a single random walk is a Brownian motion.



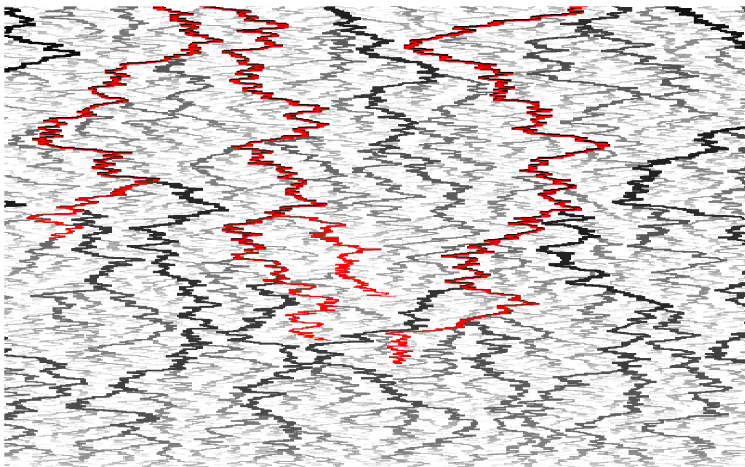
Random walk paths started at different space-time points *coalesce* as soon as they meet.



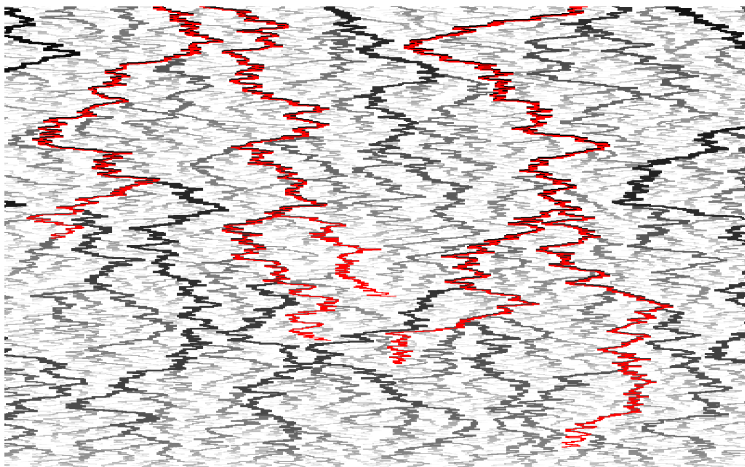
In the limit, we obtain *coalescing Brownian motions*.



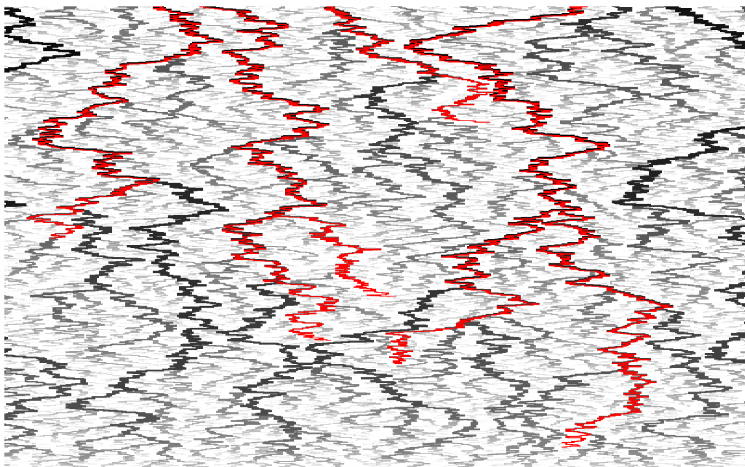
In the limit, we obtain *coalescing Brownian motions*.



In the limit, we obtain *coalescing Brownian motions*.



In the limit, we obtain *coalescing Brownian motions*.



In the limit, we obtain *coalescing Brownian motions*.



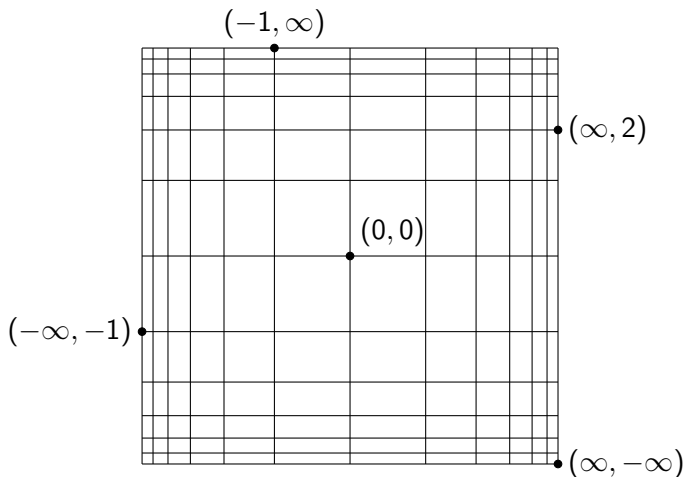
We want to give a rigorous construction of an object with the following informal description:

*Coalescing Brownian motions,  
started in each space-time point.*

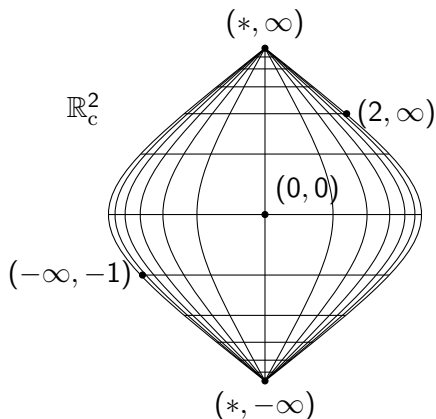
Moreover, we want to show that diffusively rescaled arrow configurations converge to such an object.

To this aim, we must first introduce the right topology.

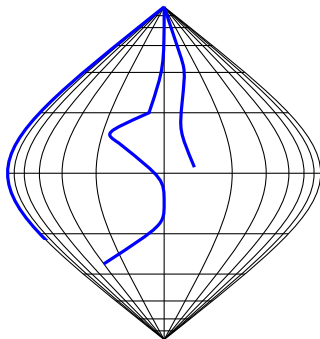
# Topological matters



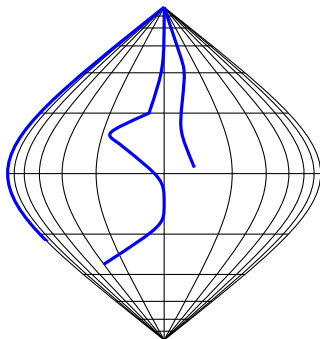
We compactify both space and time by adding points at  $\pm\infty$ .



... and then contract  $[-\infty, \infty] \times \{-\infty\}$   
and  $[-\infty, \infty] \times \{\infty\}$  to single points.

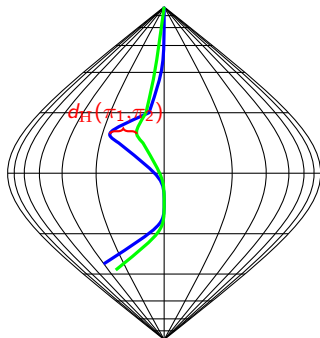


A *path* is a continuous function  $\pi : [\sigma_\pi, \infty) \rightarrow [-\infty, \infty]$ ,  
where  $\sigma_\pi \in \mathbb{R}$  is the *starting time*.



We identify a path with its graph

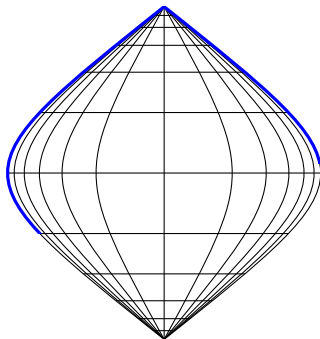
$$\{(\pi(t), t) : t \in [\sigma_\pi, \infty)\} \cup \{(*, \infty)\} \subset \mathbb{R}_c^2.$$



We equip the space  $\Pi$  of all paths with the *Hausdorff metric*

$$d_H(\pi_1, \pi_2) := \sup_{z_1 \in \pi_1} \inf_{z_2 \in \pi_2} d(z_1, z_2) \vee \sup_{z_2 \in \pi_2} \inf_{z_1 \in \pi_1} d(z_1, z_2),$$

where  $d$  is a the metric on space-time  $\mathbb{R}_c^2$ .



The set  $\mathcal{U}$  consisting of all paths in an arrow configuration plus the trivial paths that are constantly  $-\infty$  or  $+\infty$  is a compact subset of the space of all paths  $\Pi$ .

We equip the space  $\mathcal{K}(\Pi)$  of compact subsets of the path space  $\Pi$  with the Hausdorff metric:

$$d_{\text{HH}}(\mathcal{U}_1, \mathcal{U}_2) := \sup_{\pi_1 \in \mathcal{U}_1} \inf_{\pi_2 \in \mathcal{U}_2} d_{\text{H}}(\pi_1, \pi_2) \vee \sup_{\pi_2 \in \mathcal{U}_2} \inf_{\pi_1 \in \mathcal{U}_1} d_{\text{H}}(\pi_1, \pi_2).$$

Here  $d_{\text{H}}$  is the metric on  $\Pi$ .

We let  $\theta_\varepsilon : \mathbb{R}_{\text{C}}^2 \rightarrow \mathbb{R}_{\text{C}}^2$  denote the diffusive scaling map

$$\theta_\varepsilon(x, t) := (\varepsilon x, \varepsilon^2 t),$$

and set

$$\theta_\varepsilon(\pi) := \{\theta_\varepsilon(x, t) : (x, t) \in \pi\} \quad \theta_\varepsilon(\mathcal{U}) := \{\theta_\varepsilon(\pi) : \pi \in \mathcal{U}\}.$$



**Theorem [Fontes, Isopi, Newman, Ravishankar (2003)]** The set  $\mathcal{U}$  of paths in an arrow configuration (plus trivial paths) satisfies

$$\mathbb{P}[\theta_\varepsilon(\mathcal{U}) \in \cdot] \xrightarrow{\varepsilon \rightarrow 0} \mathbb{P}[\mathcal{W} \in \cdot],$$

where  $\Rightarrow$  denotes weak convergence of probability laws on  $\mathcal{K}(\Pi)$  and  $\mathcal{W}$  is a random compact set of paths called the *Brownian web*.

# The Brownian web

The Brownian web is a the unique (in law) random compact subset of paths such that:

- ▶ In each deterministic point  $z \in \mathbb{R}^2$  there almost surely starts precisely one path  $p_z \in \mathcal{W}$ .

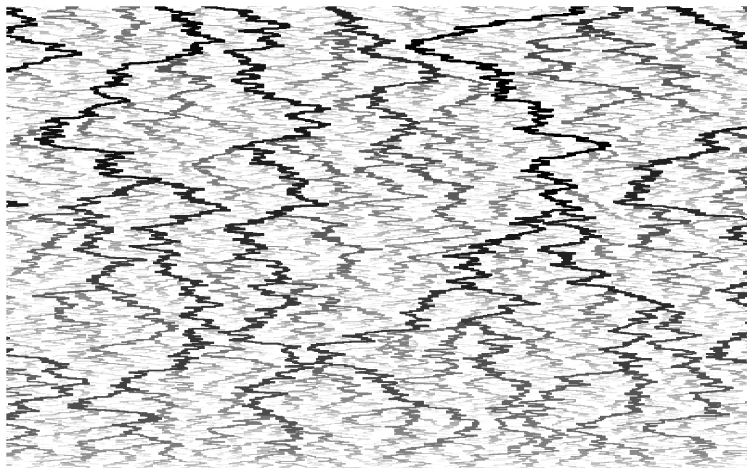
The Brownian web is a the unique (in law) random compact subset of paths such that:

- ▶ In each deterministic point  $z \in \mathbb{R}^2$  there almost surely starts precisely one path  $p_z \in \mathcal{W}$ .
- ▶ The paths  $p_{z_1}, \dots, p_{z_k}$  starting in a finite collection  $z_1, \dots, z_k \in \mathbb{R}^2$  of deterministic points are distributed as coalescing Brownian motions.

The Brownian web is a the unique (in law) random compact subset of paths such that:

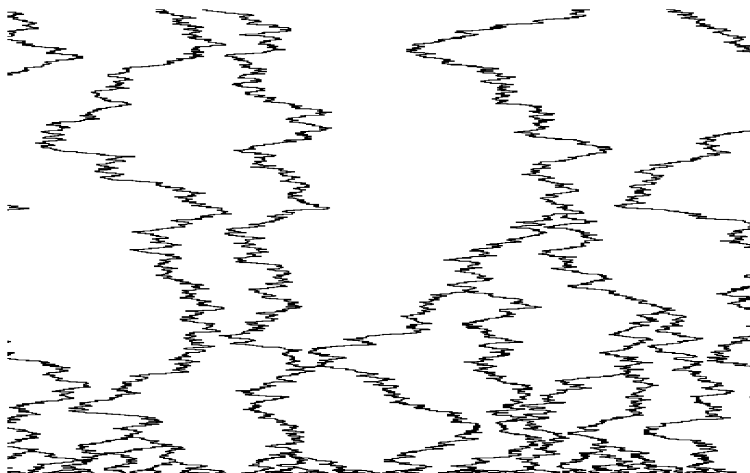
- ▶ In each deterministic point  $z \in \mathbb{R}^2$  there almost surely starts precisely one path  $p_z \in \mathcal{W}$ .
- ▶ The paths  $p_{z_1}, \dots, p_{z_k}$  starting in a finite collection  $z_1, \dots, z_k \in \mathbb{R}^2$  of deterministic points are distributed as coalescing Brownian motions.
- ▶ For each deterministic countable dense set  $\mathcal{D} \subset \mathbb{R}^d$ , the Brownian web  $\mathcal{W}$  is the closure of the set  $\mathcal{W}(\mathcal{D}) := \{p_z : z \in \mathcal{D}\}$  of paths starting in  $\mathcal{D}$ .

# The Brownian web



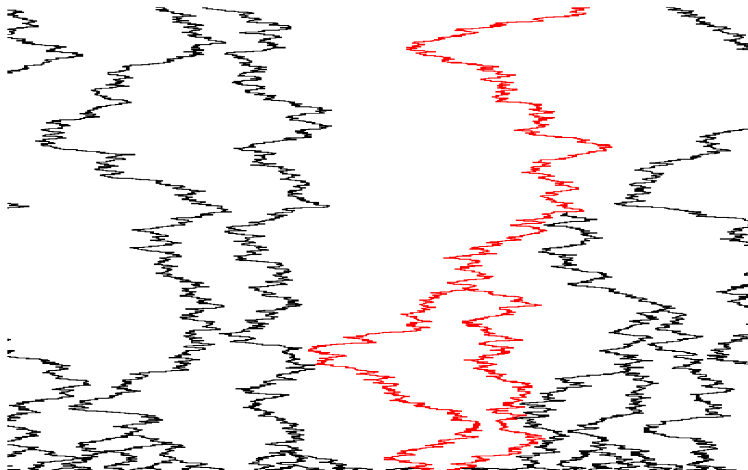
Artist's impression of the Brownian web.

# The Brownian web



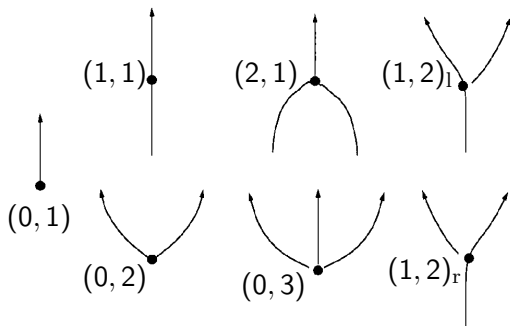
Paths starting at time zero.

# The Brownian web



Even though at deterministic points  $z$  there a.s. starts a single path  $\pi_z$ , there exist random points that are the starting point of two paths.

# Special points

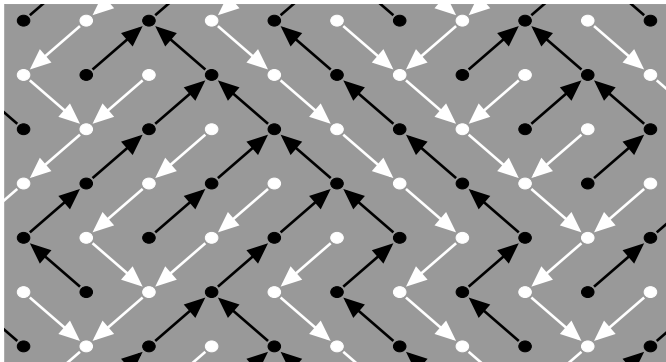


We can distinguish points in the plane according to the number of distinct paths entering and leaving a point.

In total, there are 7 types of points.

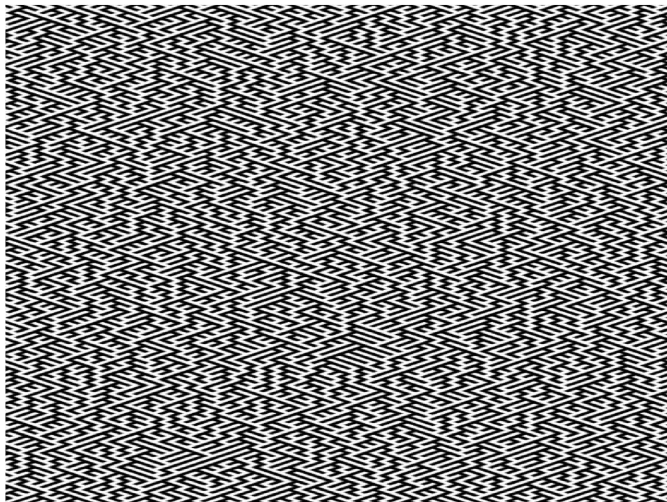


# Dual arrows



Each arrow configuration defines a *dual* arrow configuration.

# Dual Brownian web

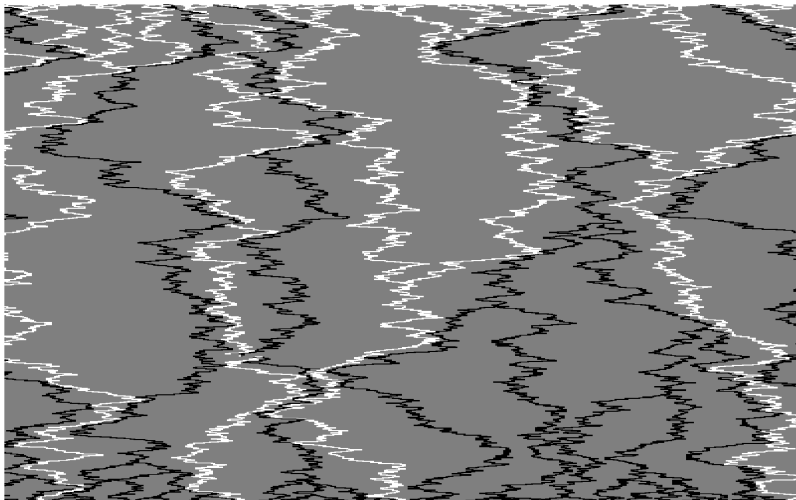


In the diffusive scaling limit, the dual arrow configuration converges to a dual Brownian web.

Associated to each Brownian web  $\mathcal{W}$ , there is a *dual* Brownian web  $\hat{\mathcal{W}}$  that is a.s. uniquely determined by  $\mathcal{W}$  and equally distributed with  $\mathcal{W}$  after a rotation over  $180^\circ$ .

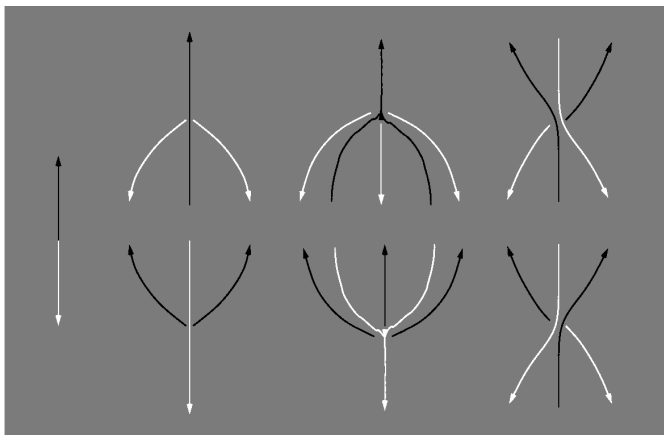
Dual paths reflect off forward paths with Skorohod reflection.

# Dual Brownian web



Forward paths (black) and dual paths (white)  
starting at two fixed times.

# Special points revisited



The structure of the dual Brownian web at special points.

# The voter model

Let  $[0, 1]^{\mathbb{Z}}$  be the space of functions  $x : \mathbb{Z} \rightarrow [0, 1]$ .

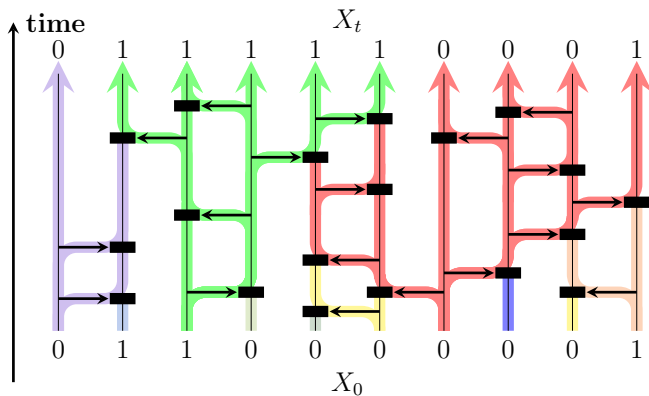
The infinite-type one-dimensional voter model is a Markov process  $(X_t)_{t \geq 0}$  with state space  $[0, 1]^{\mathbb{Z}}$ .

We call  $X_t(i) \in [0, 1]$  the *type* if site  $i \in \mathbb{Z}$  at time  $t \geq 0$ .

Initially,  $(X_0(i))_{i \in \mathbb{Z}}$  are i.i.d. uniformly distributed.

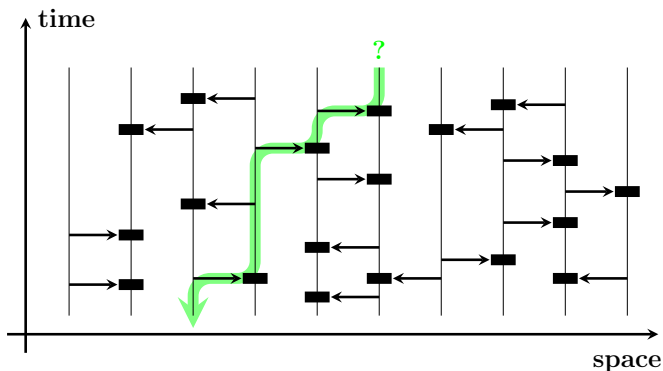
At times of a Poisson point process with intensity one, the site  $i$  selects one of its neighbours at random and copies its type.

# The voter model



A one-dimensional voter model.

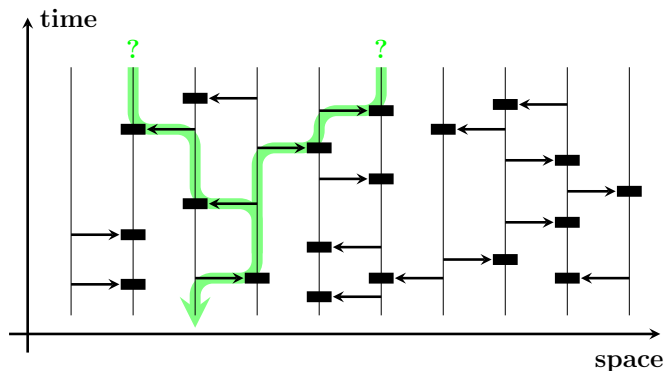
# The voter model



To find out the type of site  $i$  at time  $t$ ,  
we follow the arrows backwards.

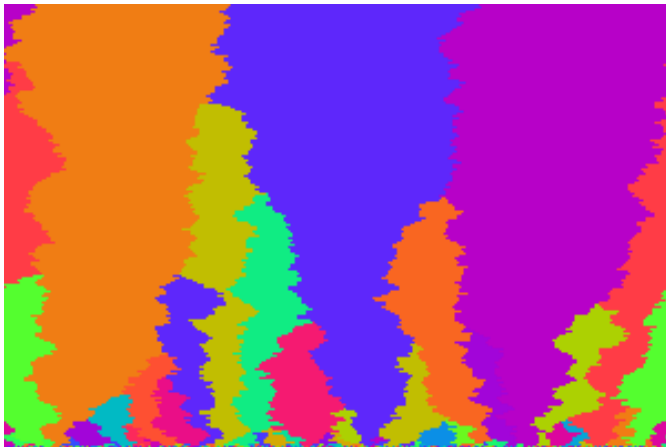


# The voter model



Ancestral lines started from different points  
coalesce when they meet.

# The voter model



In the diffusive scaling limit, the ancestral lines converge to a dual Brownian web. Arratia (1979) initiated the study of the Brownian web with the aim of describing the scaling limit of the voter model.

# The heat equation with Wright-Fisher noise

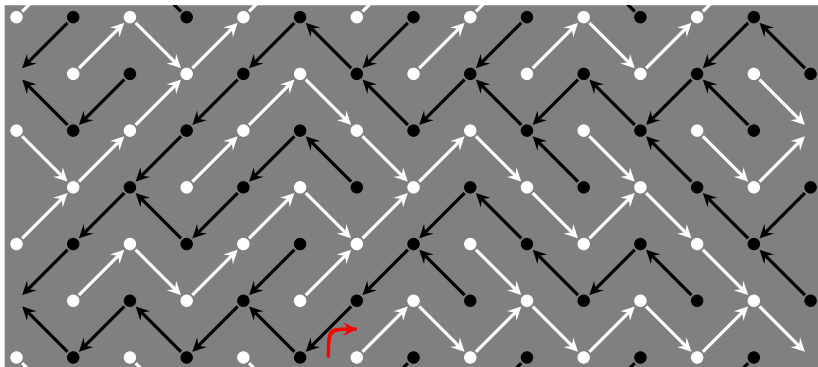
In 1988, T. Shiga studied the *heat equation with Wright-Fisher noise*

$$\frac{\partial}{\partial t} u(x, t) = \frac{1}{2} \frac{\partial^2}{\partial x^2} u(x, t) + \gamma \sqrt{u(x, t)(1 - u(x, t))} \partial W(x, t),$$

where  $W(x, t)$  is space-time white noise. He showed that solutions to this PDE can via duality be expressed in terms of systems of Brownian motions that coalesce with rate  $\gamma$  times their intersection local time.

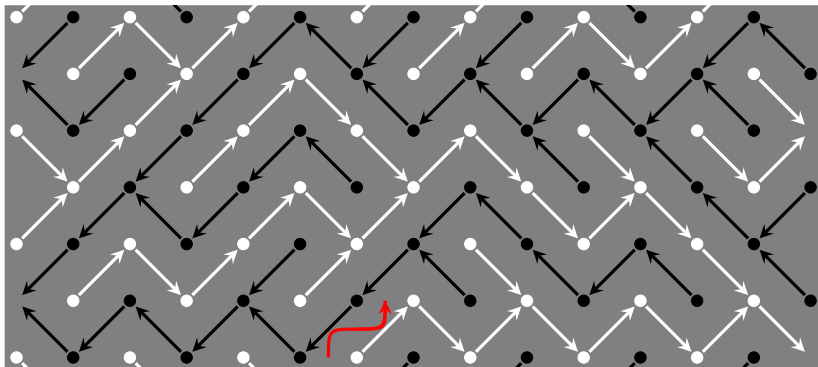
The continuum voter model with type space  $\{0, 1\}$  corresponds to the  $\gamma \rightarrow \infty$  limit of this PDE.

# The true self-repelling motion



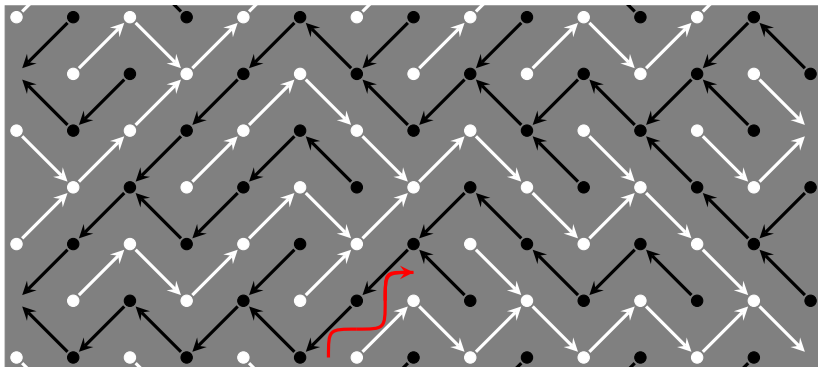
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



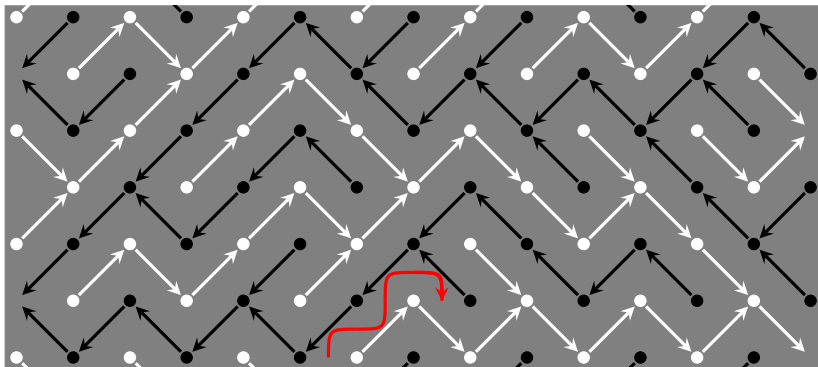
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



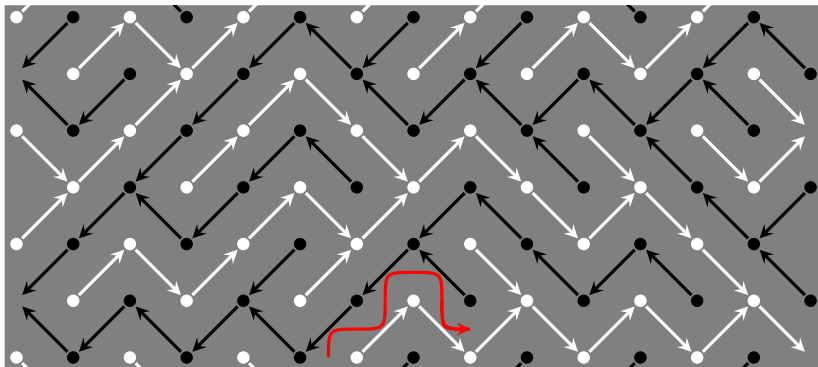
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

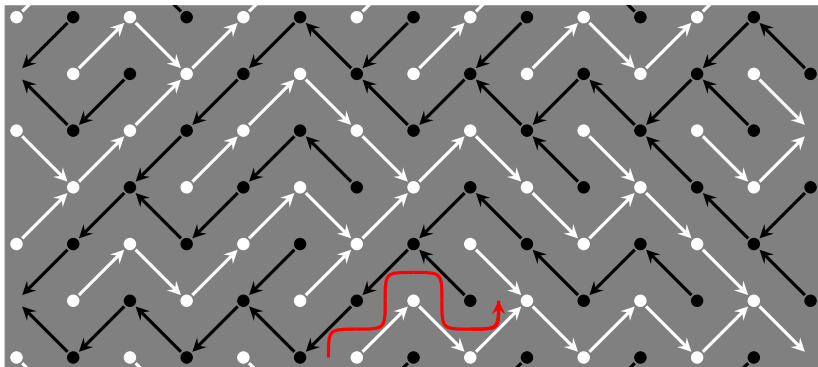
# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

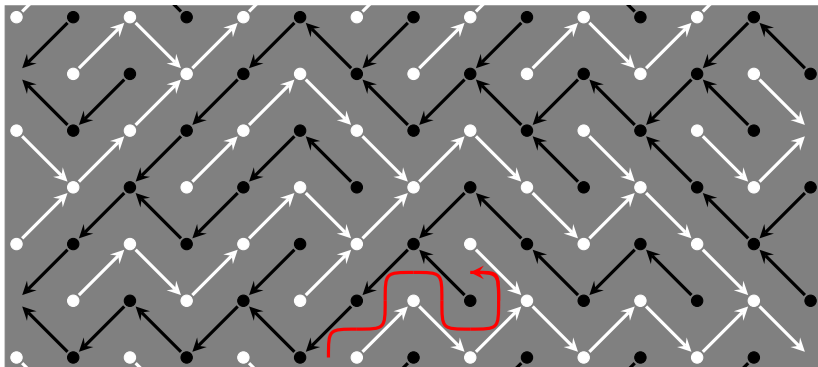


# The true self-repelling motion



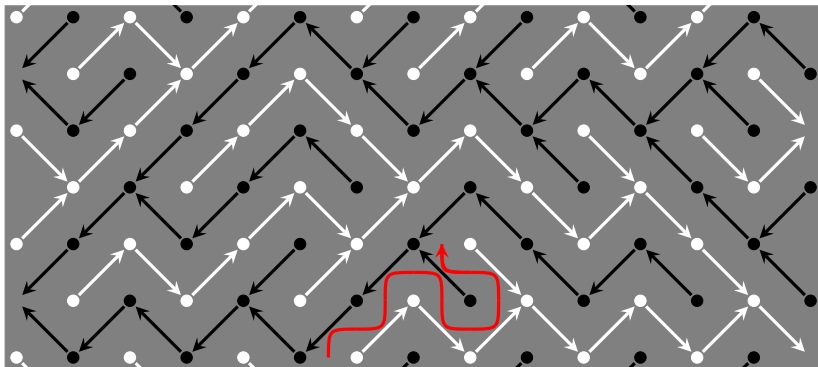
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



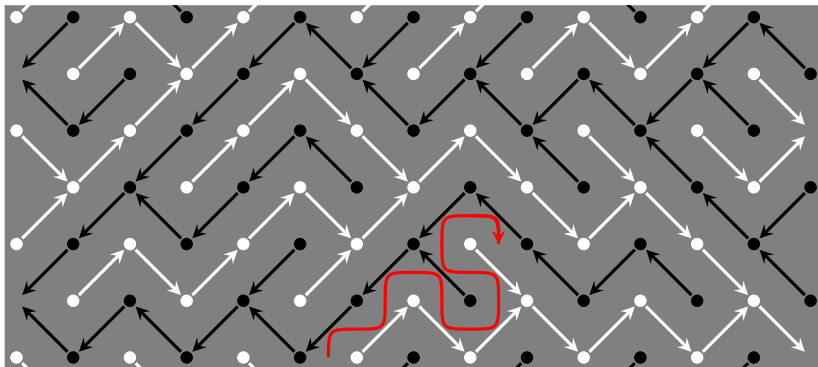
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



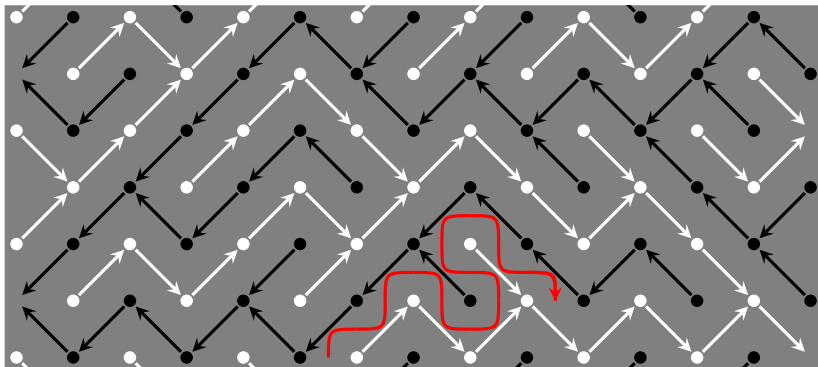
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



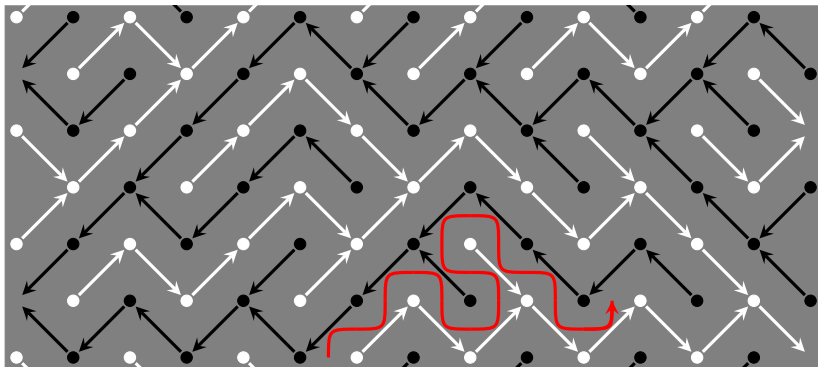
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



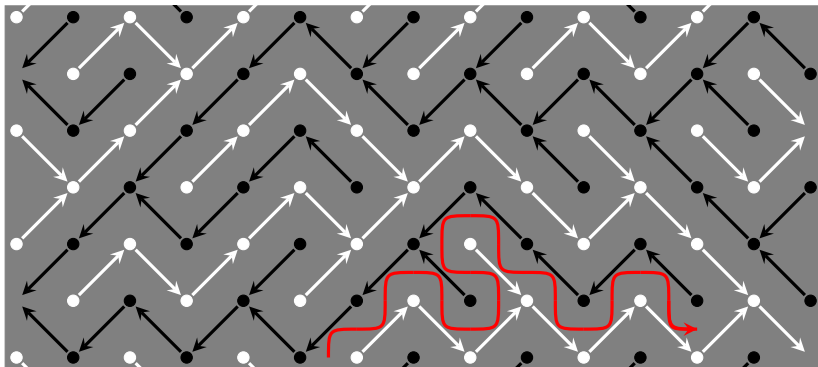
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



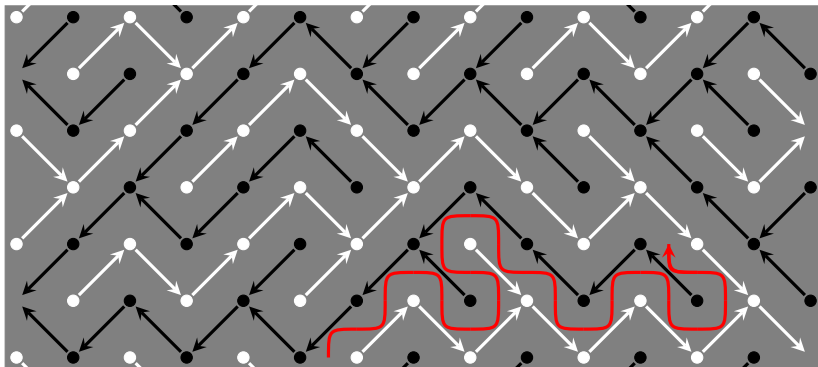
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

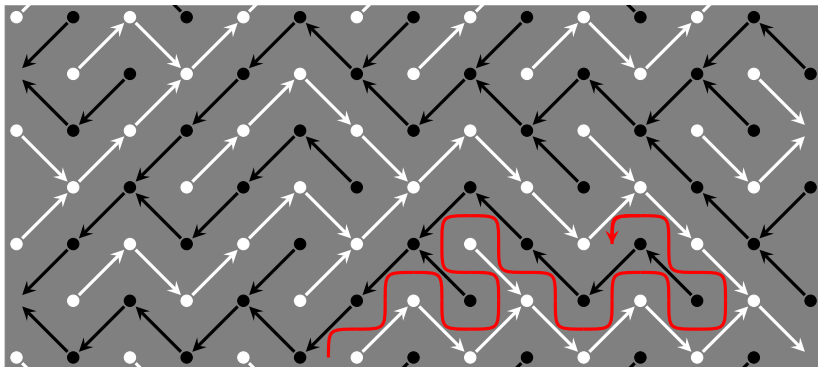
# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

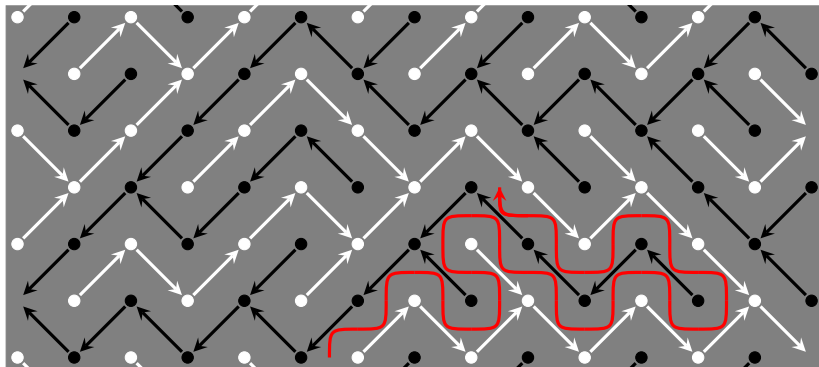


# The true self-repelling motion



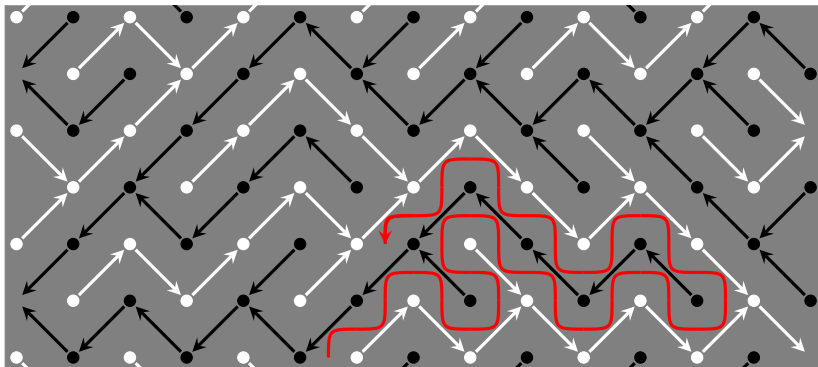
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



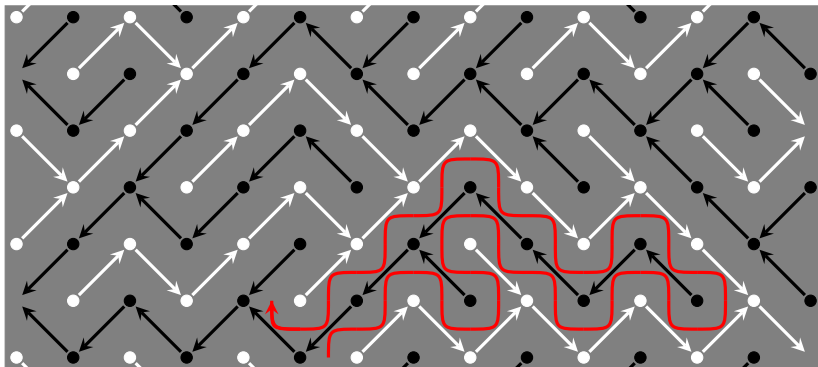
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



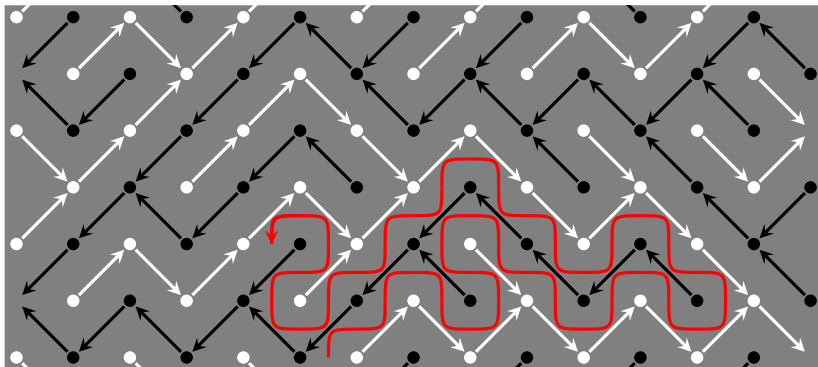
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



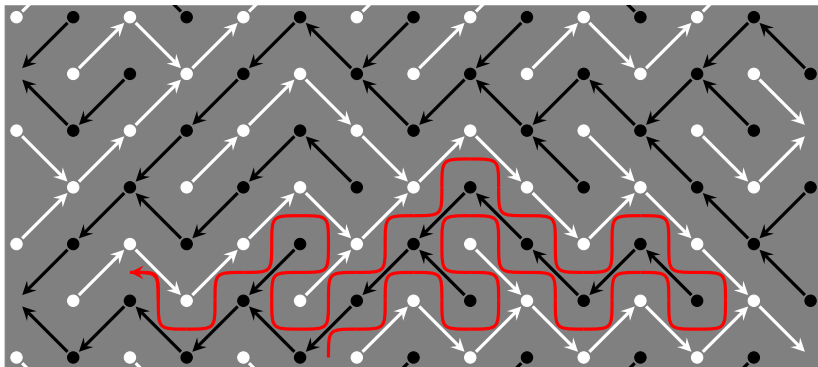
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



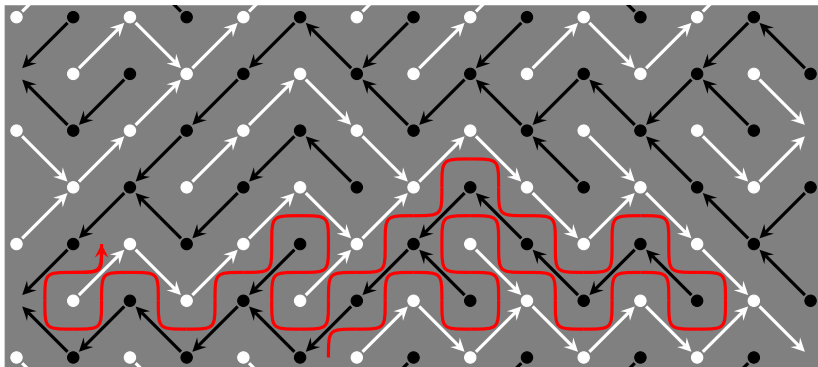
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



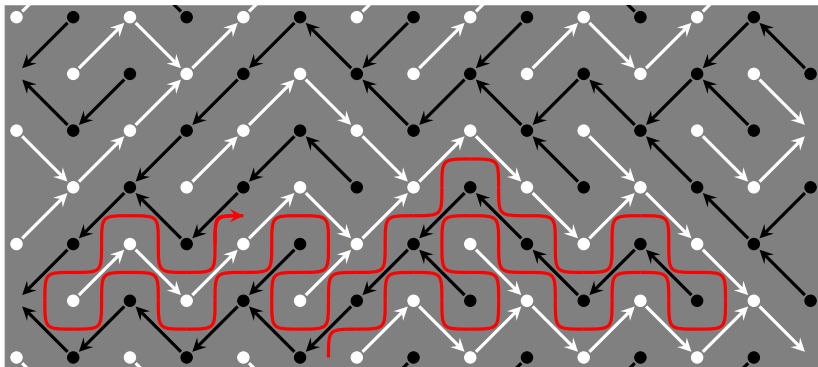
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

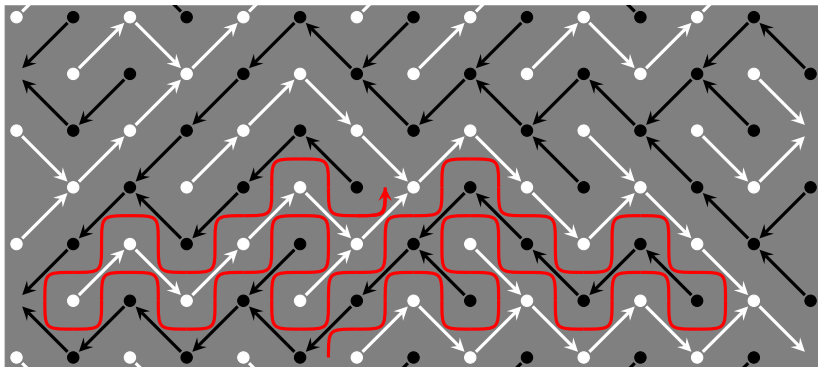
# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

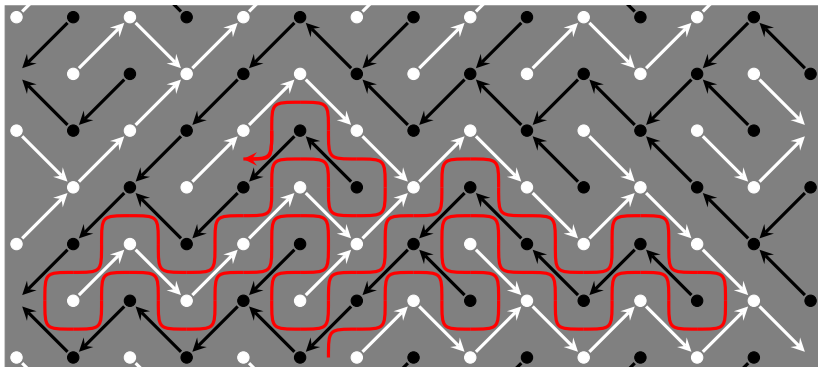


# The true self-repelling motion



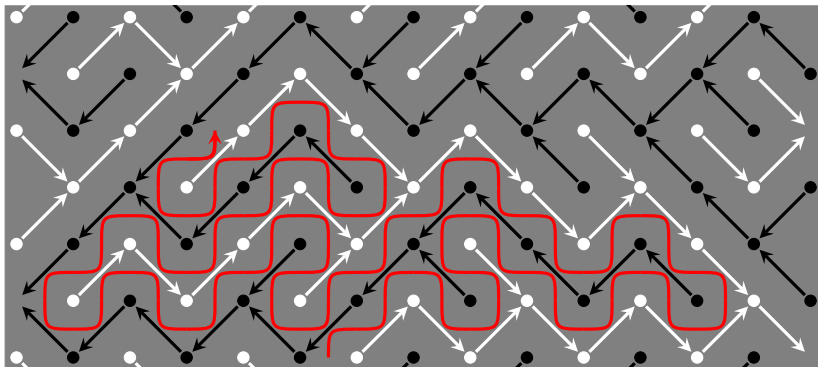
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



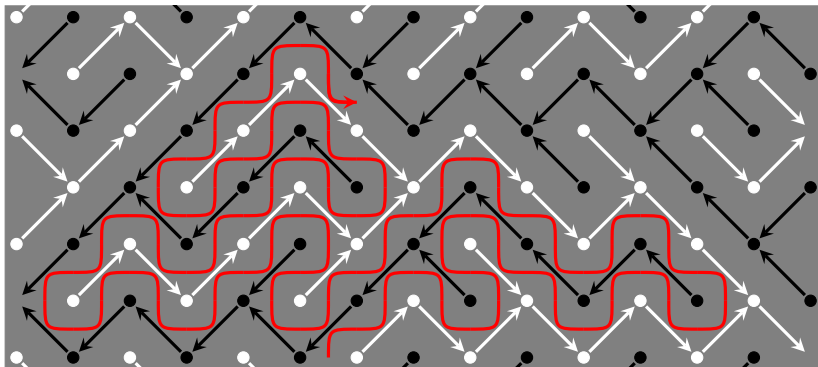
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



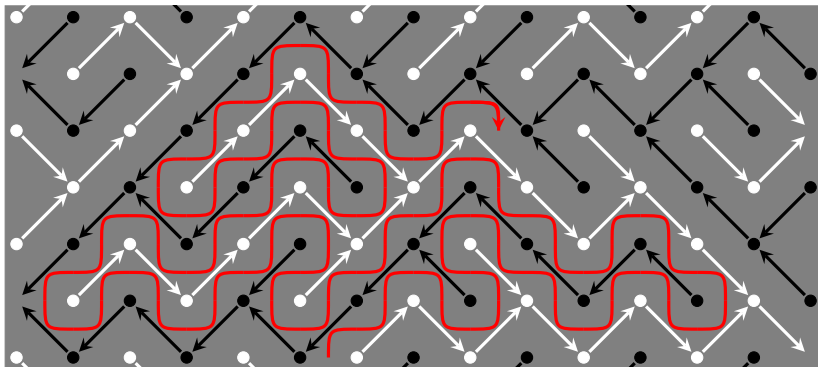
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



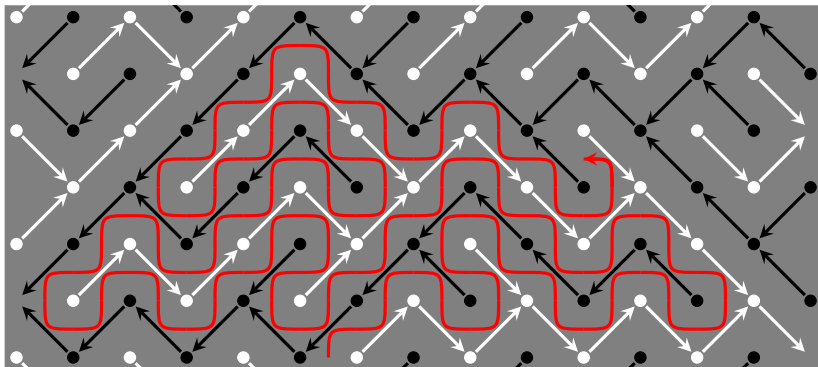
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



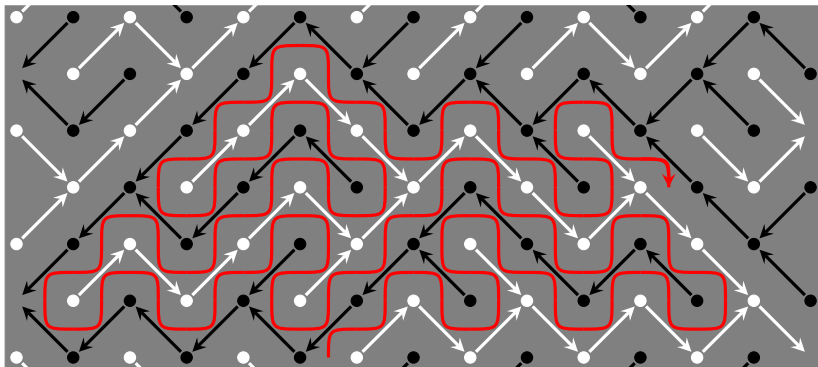
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



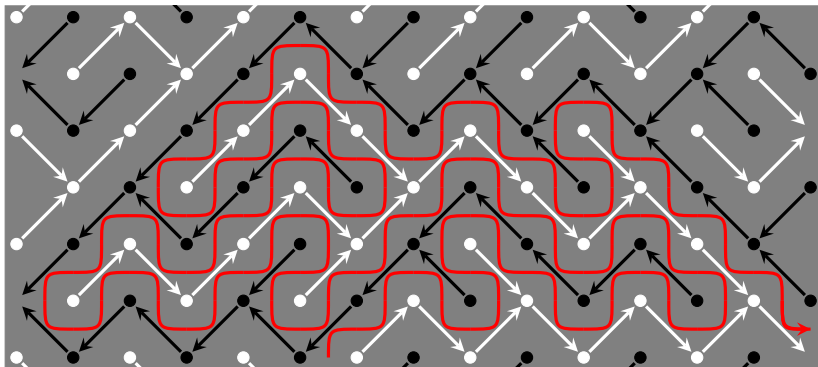
In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.

# The true self-repelling motion



In 1998, Bálint Tóth and Wendelin Werner used the Brownian web to describe the scaling limit of the *true self-repelling motion*.



# The directed spanning forest

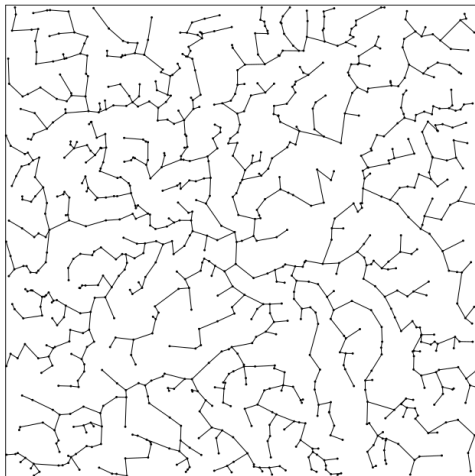
In the *Euclidean directed spanning forest*<sup>1</sup> we place points in the plane according to a Poisson point process with intensity one, and we add one extra point at the origin.

We then connect each point to the nearest point that is closer to the origin.

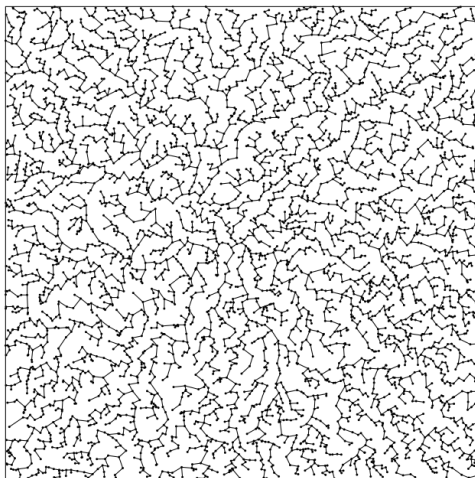
---

<sup>1</sup>Thanks to Kateřina Pawlasová for the pictures.

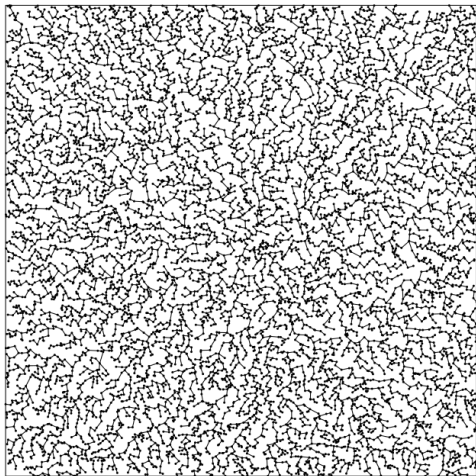
# The directed spanning forest



# The directed spanning forest



# The directed spanning forest



# The directed spanning forest

In 2021, D. Coupier, K. Saha, A. Sarkar, and V.C. Tran showed that a directed version of this model scales to the Brownian web.

# The 0-ballistic deposition model

In 2021, G. Cannizzaro and M. Hairer studied the *0-ballistic deposition model* that describes the growth of a random interface  $H_t : \mathbb{Z} \rightarrow \mathbb{N}$ .

Each site  $i \in \mathbb{Z}$  is updated with Poisson rate one, in such a way that:

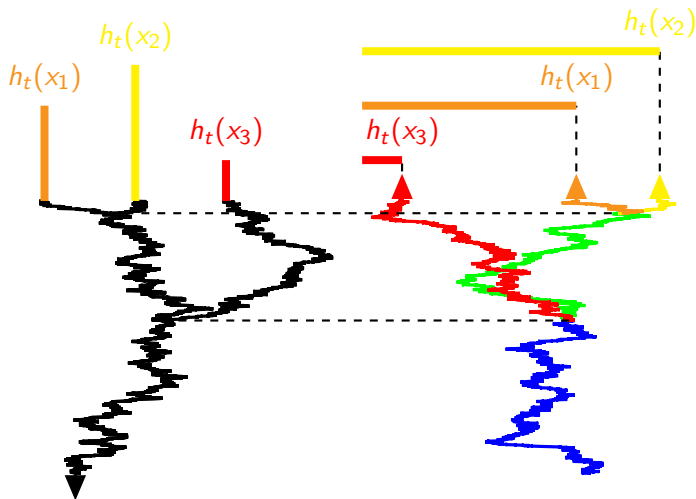
$$H_t(i) \mapsto \begin{cases} H_t(i-1) & \text{with probability } 1/3, \\ H_t(i) + 1 & \text{with probability } 1/3, \\ H_t(i+1) & \text{with probability } 1/3. \end{cases}$$

# The 0-ballistic deposition model



After subtracting the mean, we are interested in the diffusive scaling limit of the interface.

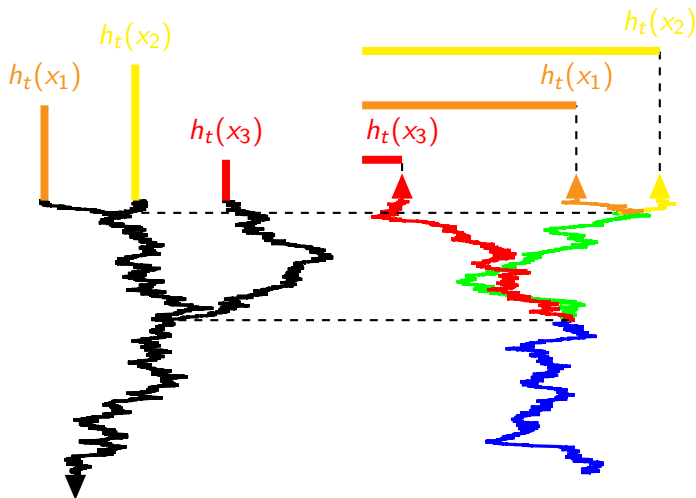
# The 0-ballistic deposition model



To determine the heights  $h_t(x_i)$  in points  $x_1, \dots, x_n$ , we first construct downward coalescing Brownian motions from  $(x_i, t)$ .

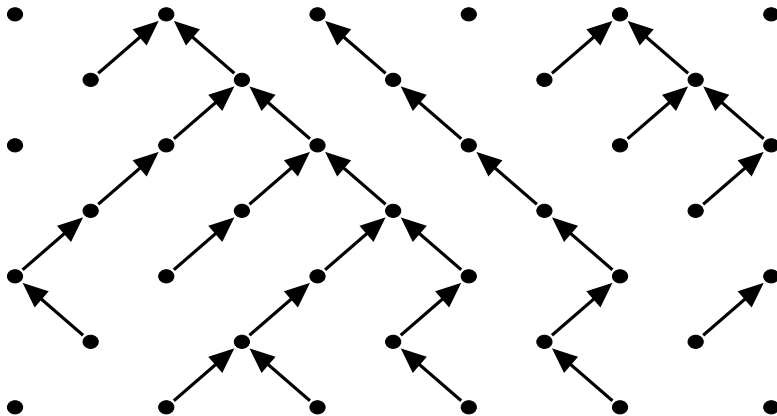


# The 0-ballistic deposition model



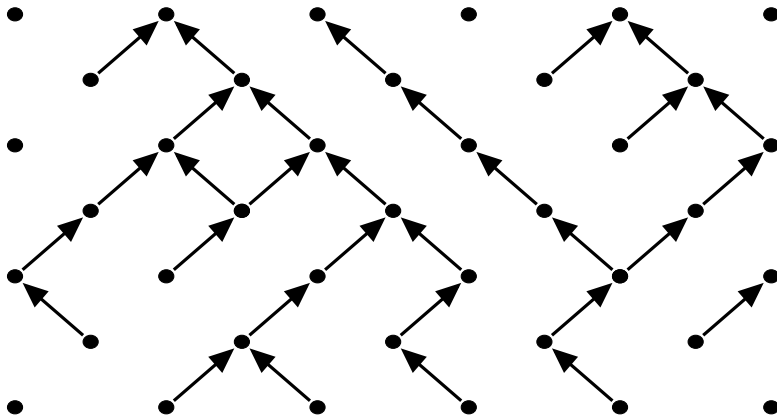
The heights are then determined by upward Brownian motions that branch with the tree structure of the downward Brownian motions.

# Branching and coalescence



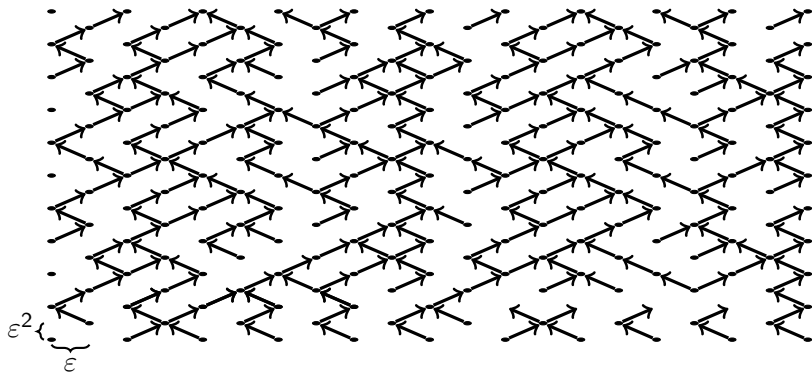
We can change an arrow configuration by drawing with probability  $\varepsilon$  two arrows, one to the left and one to the right.

# Branching and coalescence



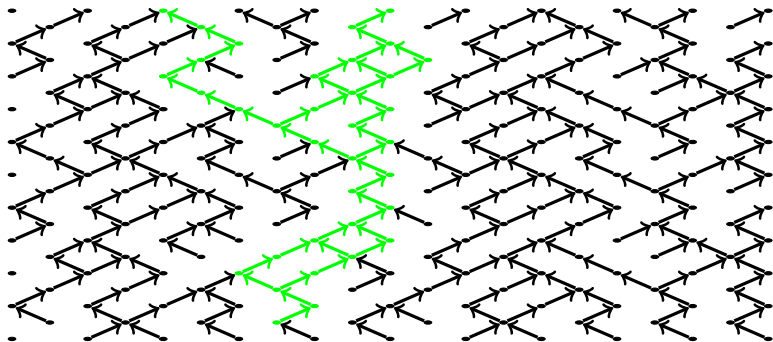
We can change an arrow configuration by drawing with probability  $\varepsilon$  two arrows, one to the left and one to the right.

# Branching and coalescence



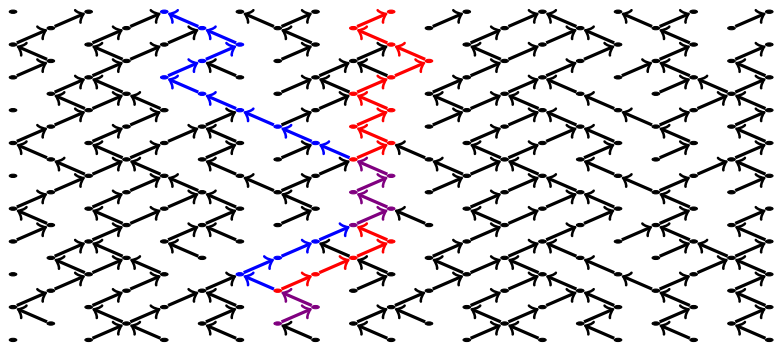
We again rescale space by  $\epsilon$  and time by  $\epsilon^2$ .

# Branching and coalescence



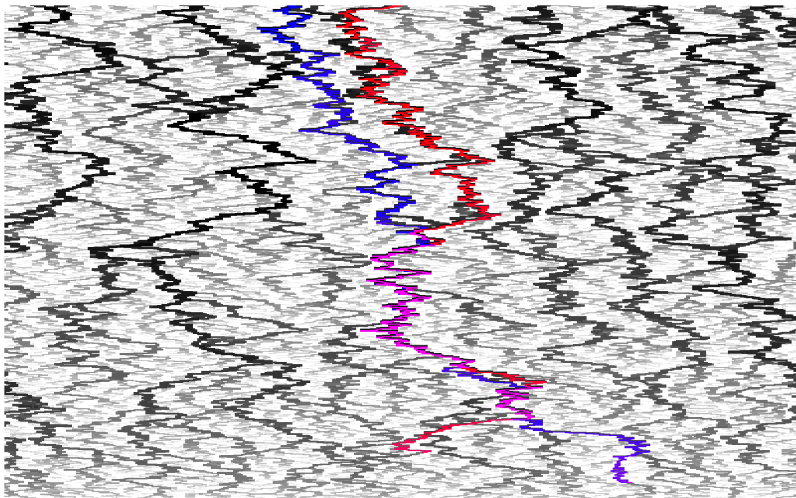
In each space-time point, there now start many different paths.

# Left and right paths



At each point, there start a unique left-most and right-most path.

# Left and right paths



After diffusive rescaling, these converge  
to drifted Brownian motions.

# Left and right paths

The joint law of left and right paths is described by the stochastic differential equation:

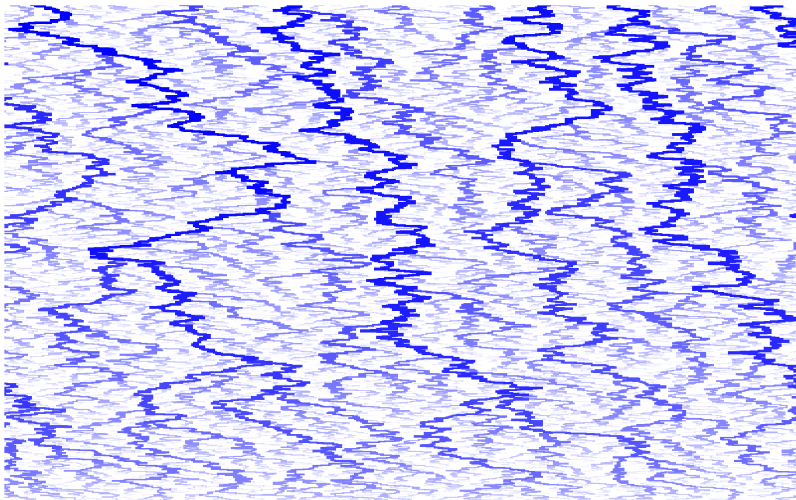
$$\begin{aligned}dL_t &= 1_{\{L_t \neq R_t\}} dB_t^l + 1_{\{L_t = R_t\}} dB_t^s - dt, \\dR_t &= 1_{\{L_t \neq R_t\}} dB_t^r + 1_{\{L_t = R_t\}} dB_t^s + dt,\end{aligned}$$

where  $B_t^l, B_t^r, B_t^s$  are independent Brownian motions, and  $L_t$  and  $R_t$  satisfy  $L_t \leq R_t$  for all  $t \geq \tau := \inf\{u \geq 0 : L_u = R_u\}$ .

The set  $\{t : L_t = R_t\}$  is nowhere dense and has positive Lebesgue measure.

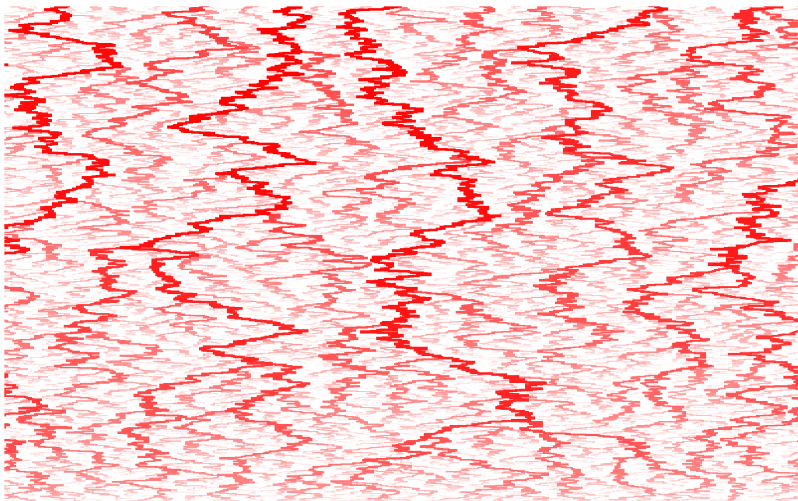


# Left Brownian web



All left-most paths form a left Brownian web...

# Right Brownian web



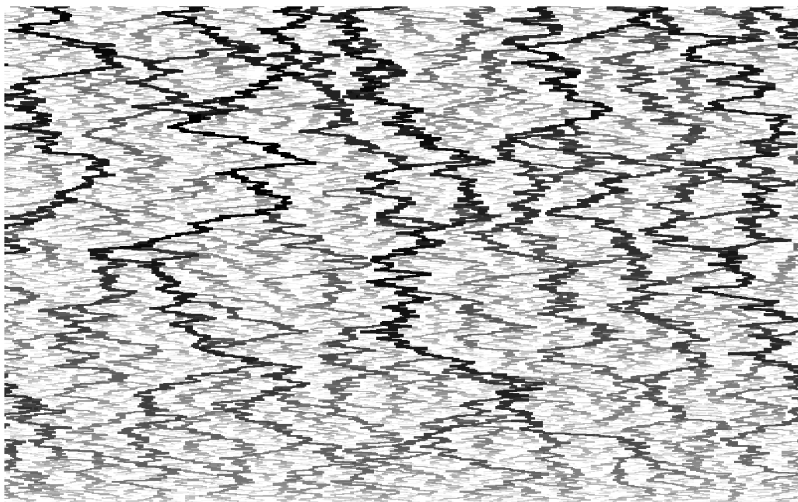
... and the right-most paths form a right Brownian web.

**Theorem [Sun, S. (2008)]** Let  $\mathcal{U}_\varepsilon$  be the collection of paths in an arrow configuration with branching probability  $\varepsilon$ . Then

$$\mathbb{P}[\theta_\varepsilon(\mathcal{U}_\varepsilon) \in \cdot] \xrightarrow[\varepsilon \rightarrow 0]{} \mathbb{P}[\mathcal{N} \in \cdot],$$

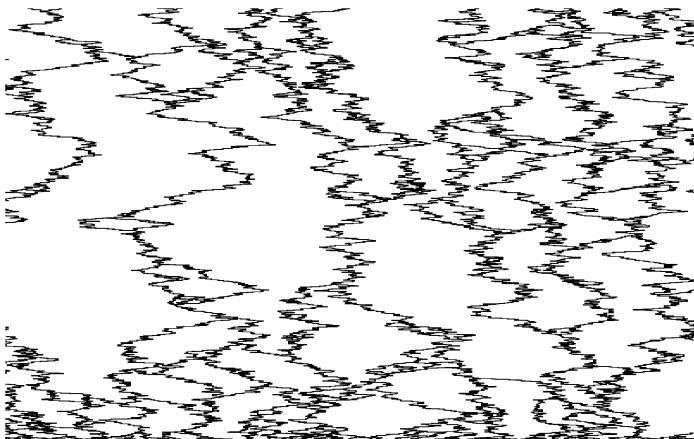
where  $\Rightarrow$  denotes weak convergence on the space  $\mathcal{K}(\Pi)$  of compact sets of paths, and the limiting object  $\mathcal{N}$  is called the *Brownian net*.

# Brownian net



Artist's impression of the Brownian net.

# Brownian net



The collection of paths starting at time zero.

# The branching-coalescing point set

For each closed set  $A \subset \mathbb{R}$ ,

$$\xi_t := \{ \pi_t : \exists \pi \in \mathcal{N} \text{ s.t. } \sigma_\pi = 0, \pi_0 \in A \} \quad (t \geq 0)$$

defines a Markov process  $(\xi_t)_{t \geq 0}$  with values in the space of closed subsets of  $\mathbb{R}$ .

- (i) Invariant law: the law of a Poisson point process with intensity 1.
- (ii) For deterministic  $t > 0$ , the set  $\xi_t$  is a.s. a locally finite subset of  $\mathbb{R}$ .
- (iii) There exists a dense set of random times at which  $\xi_t$  has no isolated points.

# Potts's model

A one-dimensional Potts model with Glauber dynamics is a Markov process  $(X_t)_{t \geq 0}$  taking values in the space  $\{1, \dots, q\}^{\mathbb{Z}}$  of functions  $x : \mathbb{Z} \rightarrow \{1, \dots, q\}$ .

For  $x \in \{1, \dots, q\}^{\mathbb{Z}}$ , let

$$N_i^x(\sigma) := \sum_{j \in \{i-1, i+1\}} \mathbf{1}_{\{x(j)=\sigma\}}$$

denote the number of neighbours of  $i \in \mathbb{Z}$  that have the value  $\sigma \in \{1, \dots, q\}$ .

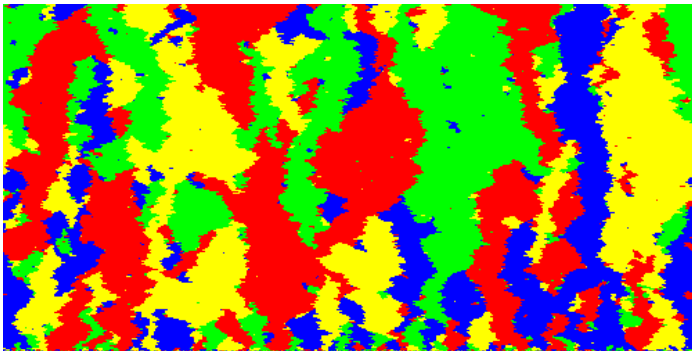
Each site  $i$  is updated with Poisson rate one and chooses a new value according to the law

$$\mu_i^x(\sigma) := \frac{1}{Z_i^x} e^{\beta N_i^x(\sigma)} \quad (\sigma \in \{1, \dots, q\}),$$

where  $Z_i^x$  is a normalisation constant.

We are interested in the low temperature limit  $\beta \rightarrow \infty$ .

# Pott's model



A low temperature one-dimensional Potts model  
with Glauber dynamics.



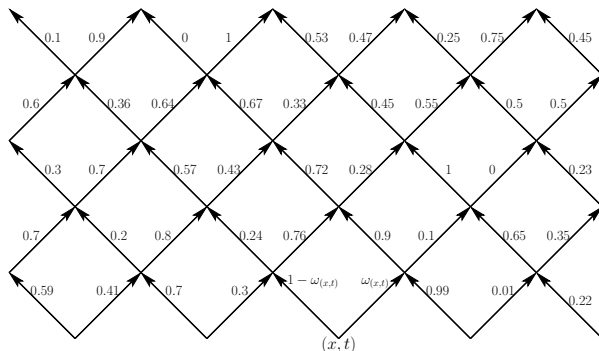
C.M. Newman, K. Ravishankar, and E. Schertzer (2015) studied coalescing random walks that branch with probability  $\varepsilon$  and die with probability  $\varepsilon^2$ .

The diffusive scaling limit is called the *Brownian net with killing*.

In 2017, they used the Brownian net with killing to describe the scaling limit of low temperature one-dimensional Potts models with Glauber dynamics.

More generally, their method applies to a wide class of voter model perturbations.

# Random space-time environment



Fix a probability law  $\mu$  on  $[0, 1]$ .

Let  $(\omega_z)_{z \in \mathbb{Z}_{\text{even}}^2}$  be i.i.d.  $[0, 1]$ -valued r.v.'s with law  $\mu$ .

# A measure-valued process

Fix some probability measure  $\rho_0$  on  $\mathbb{Z}_{\text{even}}$ , and define inductively, for  $(x, t) \in \mathbb{Z}_{\text{even}}^2$ :

$$\rho_t(x) := \omega_{(x-1, t-1)} \rho_{t-1}(x-1) + (1 - \omega_{(x+1, t-1)}) \rho_{t-1}(x+1).$$

Interpretation: in the time step from  $t$  to  $t+1$ , a  $\omega_{(x, t)}$  fraction of the mass at  $x$  is sent to  $x+1$  and the rest is sent to  $x-1$ .

Then  $(\rho_t)_{t \geq 0}$  is a Markov chain taking values alternatively in the probability measures on  $\mathbb{Z}_{\text{even}}$  and  $\mathbb{Z}_{\text{odd}}$ .

## Theorem [Le Jan & Raimond '04, Howitt & Warren '06]

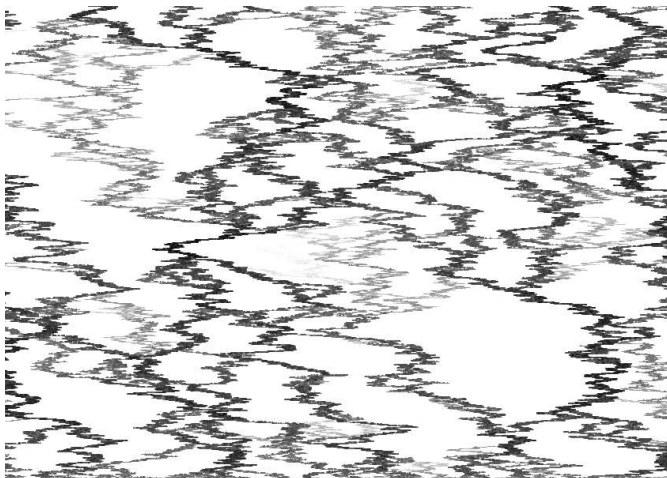
Let  $\varepsilon_n \rightarrow 0$  and let  $\rho_t^{(n)}(x)$  be Markov chains defined by splitting laws  $\mu_n$  satisfying:

$$\begin{aligned} \text{(i)} \quad & \varepsilon_n^{-1} \int 2(q - \tfrac{1}{2})\mu_n(dq) \xrightarrow[n \rightarrow \infty]{} \beta, \\ \text{(ii)} \quad & \varepsilon_n^{-1} \int q(1 - q)\mu_n(dq) \xrightarrow[n \rightarrow \infty]{} \nu(dq), \end{aligned}$$

with  $\beta \in \mathbb{R}$  and  $\nu$  a finite measure on  $[0, 1]$ .

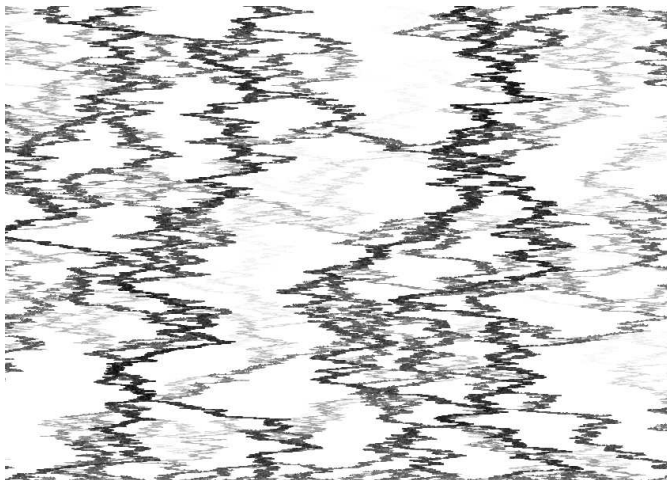
Rescale diffusively:  $\tilde{\rho}_{\varepsilon_n^2 t}^{(n)}(\varepsilon_n x) := \rho_t^{(n)}(x)$ . Then  $\tilde{\rho}^{(n)} \Rightarrow \rho$ , where  $(\rho_t)_{t \geq 0}$  is a Markov process taking values in the probability measures on  $\mathbb{R}$ , with dynamics characterized by  $\beta$  and  $\nu$ .

# Howitt-Warren flows



The equal splitting flow:  $\beta = 0$  and  $\nu = \delta_{1/2}$ .

# Howitt-Warren flows



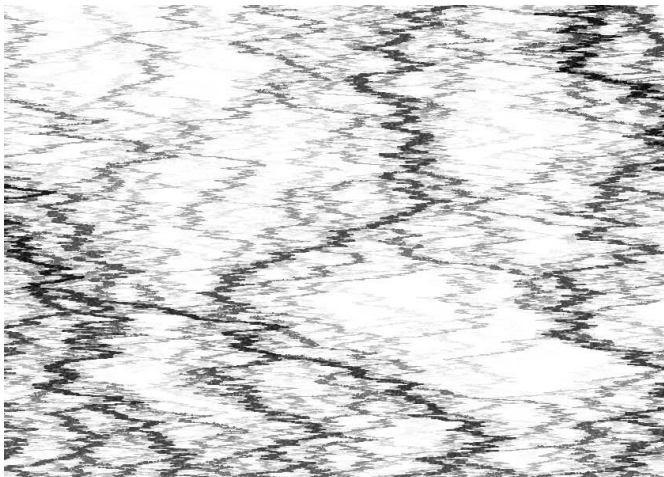
The process with  $\beta = 0$  and  $\nu(dq) = 6q(1 - q)dq$ .

# Howitt-Warren flows



Le Jan-Raimond flow:  $\beta = 0$  and  $\nu(dq) = dq$  (reversible!).

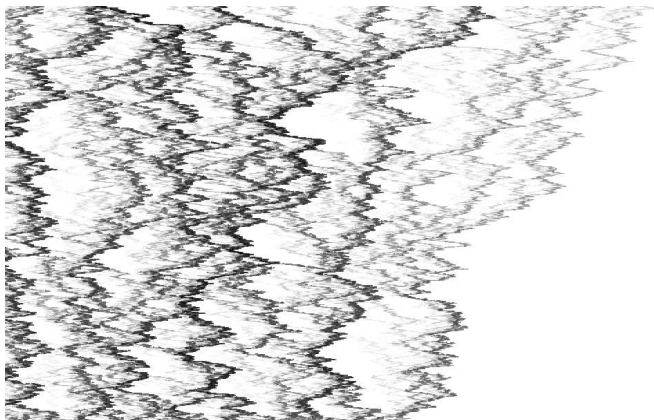
# Howitt-Warren flows



The erosion flow:  $\beta = 0$  and  $\nu = \frac{1}{2}\delta_0 + \frac{1}{2}\delta_1$ .



# Howitt-Warren flows



One-sided erosion flow:  $\beta = 0$  and  $\nu = \delta_1$ .

In 2010, E. Schertzer, R. Sun & J.S. showed that Howitt-Warren flows can be constructed with the help of the Brownian web.

If the speeds

$$\beta_+ := \beta + 2 \int q^{-1} \nu(dq),$$
$$\beta_- := \beta - 2 \int (1 - q)^{-1} \nu(dq),$$

are finite, then they are supported on a Brownian net with left speed  $\beta_-$  and right speed  $\beta_+$ .